

Review: $(M = T^*X, \omega, H)$
 (symplectic manifold) $\xleftarrow{\text{configuration space}}$ $\xleftarrow{\text{Ham. f}^n \text{ in } C^\infty(M, \mathbb{R})}$

eg.: $H \in C^\infty(M) \supseteq f^{\text{ns}}$ polynomial in cotangent fibres
 $\bigoplus_{k=0}^{\infty} \Gamma(X, \text{Sym}^k TX)$ $\xleftarrow{\text{symmetric multivectors}}$

$k=0$) $V \in C^\infty(X, \mathbb{R})$ "potential"

1) $Y \in \mathfrak{X}(X)$

2) $T_g = \frac{1}{2} f_{g^{-1}}$ g Riem. metric "kinetic energy"

Most important example: 0) + 2) :

$$H = T_g + V \quad V \in C^\infty(X, \mathbb{R})$$

$$X_H = \underbrace{X_g}_{\substack{\uparrow \\ \text{geodesic} \\ \text{v. field}}} + \underbrace{X_V}_{\substack{\uparrow \\ \text{deflection} \\ \text{due to force gen. by } V.}}$$

Recall: Hamilton's eqⁿ of motion:

$$\begin{aligned} \dot{x}^i &= \{x^i, H\} \\ \dot{p}_i &= \{p_i, H\} \end{aligned}$$

flow
of
 X_H

for X_g :

$$\textcircled{1} \begin{cases} \dot{x} = g^{-1} p \\ \dot{p}_i = \frac{1}{2} (\partial_k g^{ij}) p_i p_j \end{cases}$$

sol^{ns} are
geodesics

$$\text{accel.} = \nabla_{\frac{\partial}{\partial t}} \dot{x} = 0$$

$$\text{for } X_V = -\omega^{-1}(dV) = - (dp_i \wedge dx^i)^{-1} (\partial_k V dx^k)$$

$$\Rightarrow \textcircled{2} \begin{cases} \dot{x} = 0 \\ \dot{p}_i = -\frac{\partial V}{\partial x^i} \end{cases}$$

combining: $\nabla_{\frac{\partial}{\partial t}} \dot{x} = - \underbrace{g^{-1}(dV)}_{\text{grad } V = \nabla V} = g^{-1} F$

$$F = -dV$$

$$\in \Omega^1(X)$$

Force assoc.
to potential V .

$$F = g a = g \nabla_{\partial_t} \dot{x}$$

g plays role of mass.

Ex.: Simple harmonic oscillator



equip X w/ Riem. metric $g = m dx \otimes dx$

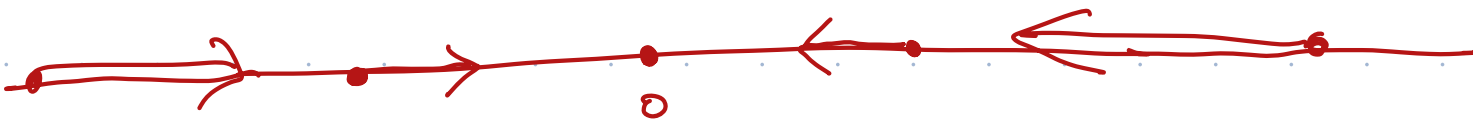
$$\Rightarrow T_g = \frac{1}{2} m^{-1} p^2$$

provide law of force for a spring (Hooke's law)

$$V = \frac{1}{2} k x^2 \quad k \in \mathbb{R}_{>0}$$



$$\begin{aligned}
 -\text{grad} V &= -g^{-1}(dV = kx dx) \\
 &= -kx \frac{\partial}{\partial x}
 \end{aligned}$$



$$H_{\text{SHO}} = \frac{1}{2} m^{-1} p^2 + \frac{1}{2} k x^2$$

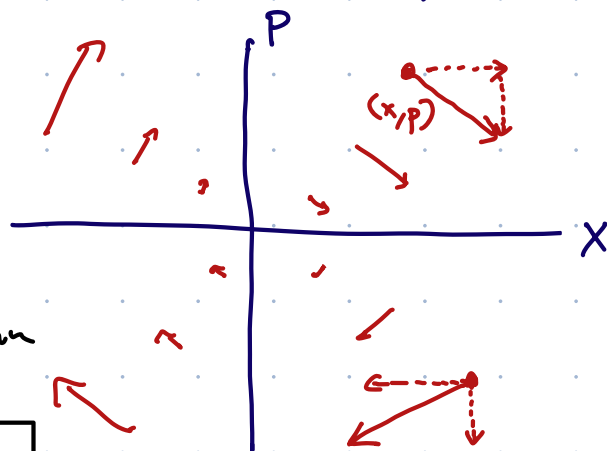
Ham. system for $X = (\mathbb{R}, x)$
 $T^*X \cong (\mathbb{R} \times \mathbb{R}, (x, p))$

$$X_H = m^{-1} p \frac{\partial}{\partial x} - kx \frac{\partial}{\partial p}$$

special property: X_H linear!

$\Rightarrow$ express Hamilton's eqns in matrix form

$$\frac{d}{dt} \begin{pmatrix} x \\ p \end{pmatrix} = \underbrace{\begin{pmatrix} 0 & m^{-1} \\ -k & 0 \end{pmatrix}}_A \begin{pmatrix} x \\ p \end{pmatrix}$$



solⁿ: $\begin{pmatrix} x(t) \\ p(t) \end{pmatrix} = e^{tA} \begin{pmatrix} x(0) \\ p(0) \end{pmatrix}$

$$A^2 = -\frac{k}{m} I = -\omega^2 I \quad \underbrace{(\omega^{-1} A)}_J^2 = -I$$

"operator behaves like i on $\mathbb{C}$ "

$$J = \left(\sqrt{\frac{k}{m}} \right)^{-1} \begin{pmatrix} 0 & m^{-1} \\ -k & 0 \end{pmatrix} = \begin{pmatrix} 0 & \frac{1}{\sqrt{km}} \\ -\sqrt{km} & 0 \end{pmatrix}$$

$$\begin{pmatrix} x(t) \\ p(t) \end{pmatrix} = \underbrace{e^{t\omega J}} \begin{pmatrix} x(0) \\ p(0) \end{pmatrix} = \left(\cos(\omega t) I + \sin(\omega t) J \right) \begin{pmatrix} x(0) \\ p(0) \end{pmatrix}$$

$$e^{i\theta} = \cos\theta + i\sin\theta$$

1st component: $x(t) = \cos(\omega t) x(0) + \frac{1}{\sqrt{km}} \sin(\omega t) p(0)$

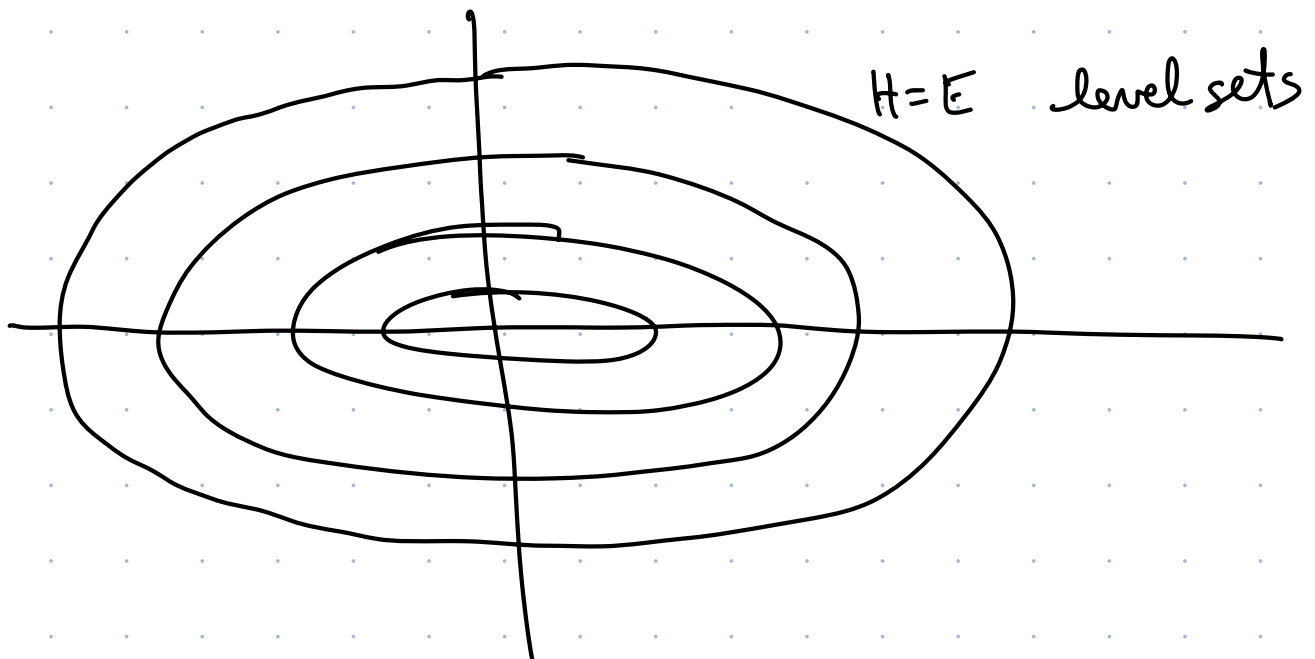
$p = m\dot{x}$

$$x(t) = \cos(\omega t) x(0) + \omega^{-1} \sin(\omega t) \dot{x}(0)$$

sinusoidal oscill. at freq ω

trajectories coincide w/ level sets of H :

$$\frac{1}{2}m^{-1}p^2 + \frac{1}{2}kx^2 = E \quad \text{constant for each solution}$$



What are the possible energies? $H \in \mathbb{R}_{\geq 0}$

whereas frequency is $\omega = \sqrt{\frac{k}{m}}$ fixed.

notice : $H = \frac{1}{2} m^{-1} p^2 + \frac{1}{2} k x^2$

we could do a $\underbrace{\text{canonical transformation}}_{\text{symplectomorphism}}$ change of variables:

$$\tilde{x} = \sqrt{k} x$$

$$\tilde{p} = \frac{1}{\sqrt{k}} p$$

$$H = \frac{1}{2} m^{-1} (k \tilde{p}^2) + \frac{1}{2} k \frac{\tilde{x}^2}{k}$$

$$= \frac{1}{2} \left(\frac{k}{m} \right) \tilde{p}^2 + \frac{1}{2} \tilde{x}^2$$

$$\tilde{x} = c x$$

$$\tilde{p} = c^{-1} p$$

$$H = \frac{1}{2} \underbrace{m^{-1} c^2} p^2 + \frac{1}{2} \underbrace{k c^{-2}} x^2$$

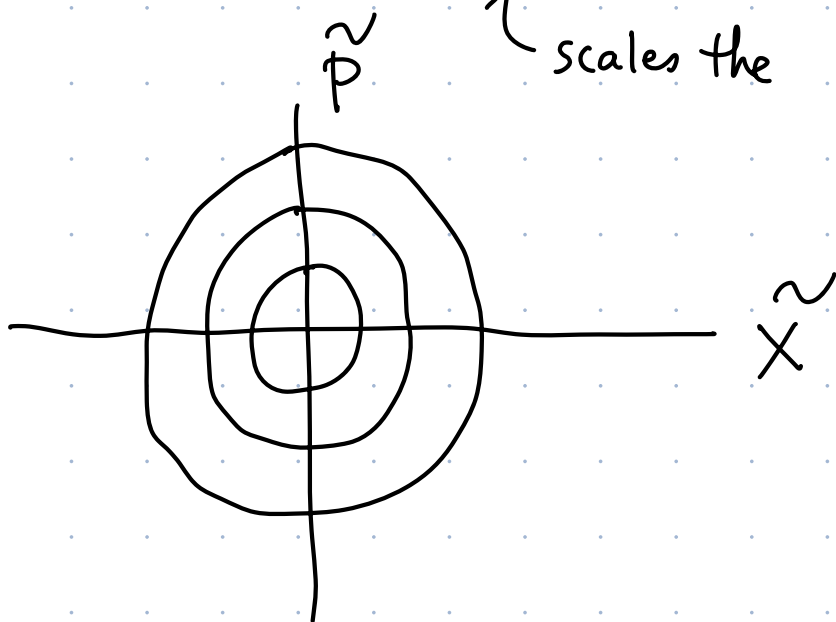
$$m^{-1} c^2 = k c^{-2}$$

$$c^4 = k m$$

$$H = \frac{1}{2} \sqrt{\frac{k}{m}} (\tilde{p}^2 + \tilde{x}^2)$$

$$\frac{c^2}{m} = \frac{\sqrt{k m}}{m} = \sqrt{\frac{k}{m}}$$

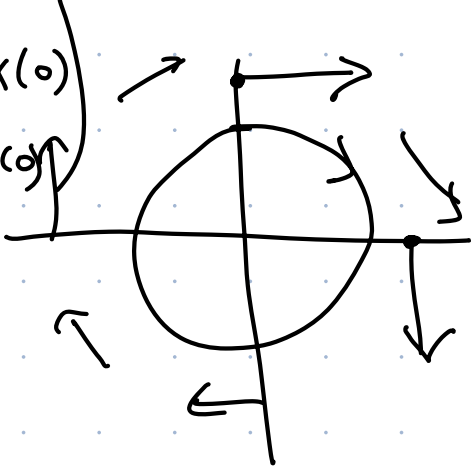
↑ scales the Hamiltonian + v. field



Standard SHO $\hbar = m = 1$ $H = \frac{1}{2} (p^2 + x^2)$

has v. field $X_H = p \frac{\partial}{\partial x} - x \frac{\partial}{\partial p}$ (standard rotation in plane)

flow : $\begin{pmatrix} \cos(t) & \sin(t) \\ -\sin(t) & \cos(t) \end{pmatrix} \begin{pmatrix} x(0) \\ p(0) \end{pmatrix}$



ω enters as a scale factor
determining angular speed.

This system has remarkable property :

- all its trajectories are closed (i.e. periodic)
- the periods all agree and are given by

period $\boxed{T = \frac{2\pi}{\omega}}$ ← transcendental number!

How is this period computed? $\int_{\text{one cycle}} dt = T$

$$\frac{dx}{dt} = \{x, H\}$$

$$\frac{dp}{dt} = \{p, H\}$$

$H^{-1}(E)$ fix energy



$$\int_{H^{-1}(E)} \left(dt = \frac{dx}{\{x, H\}} \right)$$

if we use coordinate x
on curve $\mathcal{C}$

$$T = \int_{\mathcal{C}} \frac{dx}{\{x, H\}}$$

$$T = \int_C \frac{dx}{\sqrt{2E - x^2}}$$

$$\frac{1}{2}(x^2 + p^2) = E$$

$$p = \sqrt{2E - x^2} \quad \left\{ x, \frac{1}{2}(p^2 + x^2) \right\}$$

$$2 \cdot \frac{1}{2} p \{x, p\}$$

$$\{x, p\} = \frac{\partial x}{\partial x} \cdot \frac{\partial p}{\partial p} - \cancel{\frac{\partial x}{\partial p}} \cancel{\frac{\partial p}{\partial x}} = 1$$

$$\{x, x\} = 0$$

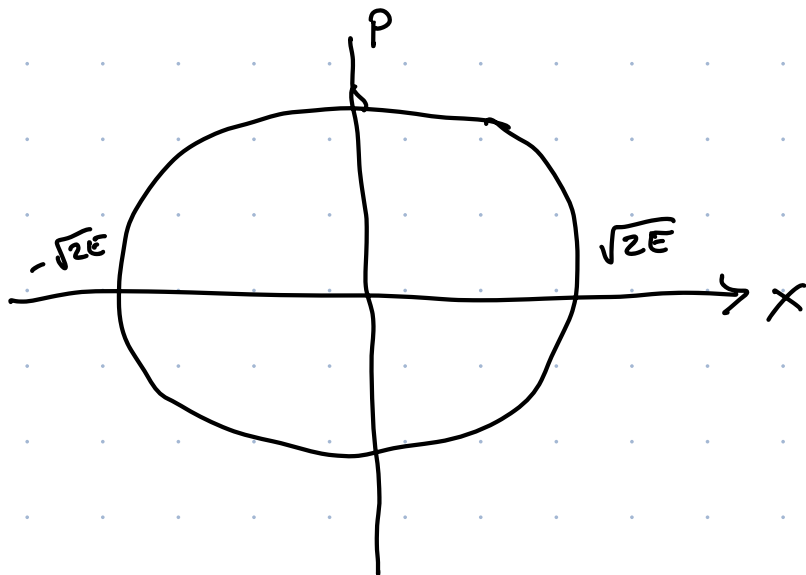
$$T = \int_{-\sqrt{2E}}^{\sqrt{2E}} \frac{dx}{\sqrt{2E - x^2}}$$

positive sq. rt.

$$+ \int_{\sqrt{2E}}^{-\sqrt{2E}} \frac{dx}{\underbrace{-\sqrt{2E - x^2}}_{\text{pos sq. rt.}}}$$

$$= 2 \int_{-\sqrt{2E}}^{\sqrt{2E}} \frac{dx}{\sqrt{2E - x^2}} = 2\pi$$

Algebraic
geometry
"periods"

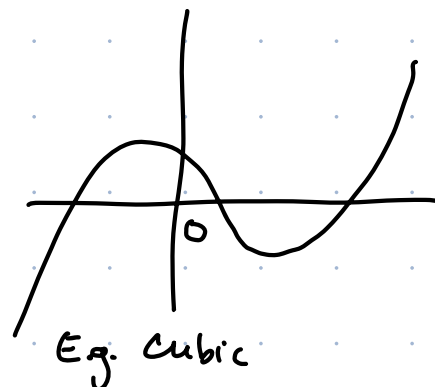


Eg 2:

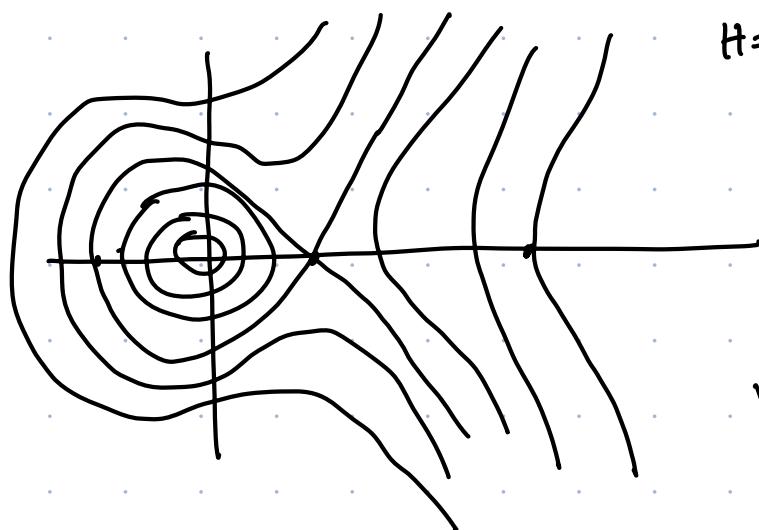
$$X = \mathbb{R}$$

$$g = g_{\text{Euc}} = dx \otimes dx$$

$$V = \text{polynomial in } x$$



$$H = \frac{1}{2} p^2 + V$$



$H = \text{const.}$

$H = 0$

V cubic $\Rightarrow$ family of elliptic curves

$$T = \text{period} = \oint_{\text{trajectory}} \frac{dx}{\sqrt{2E - V}}$$

$\nwarrow$ elliptic integral

if $\deg V > 3$

"hyperelliptic integral"

$$\dim X = 2$$

$$X = \mathbb{R}^2 \quad x^1, x^2$$

$$M = T^*X = \mathbb{R}^2 \times \mathbb{R}^2 = \mathbb{R}^4$$

4-dimensional phase space

w/ coords p_1, p_2, x^1, x^2

$$H = H(p_1, p_2, x^1, x^2)$$

$$= \frac{1}{2} (p_1^2 + p_2^2 + (x^1)^2 + (x^2)^2)$$

eg.

$$\omega = dp_1 \wedge dx^1 + dp_2 \wedge dx^2$$

(x^1, x^2 are any coords on S^2)

$$\text{Level set: } H^{-1}(E) \subset T^*X$$

E const.

$$\left\{ (x^1, x^2, p_1, p_2) : (x^1)^2 + (x^2)^2 + p_1^2 + p_2^2 = 2E \right\}$$

$$S^3 \subset \mathbb{R}^4$$

3d

sphere.

$\Rightarrow$ trajectory lies on $S^3(\sqrt{2E})$

3-sphere of radius $\sqrt{2E}$