

Geometry of Quantum Mechanics

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Hamiltonian Mechanics :

configuration space X
 n coordinates (q^1, \dots, q^n)

a covector at $x \in X$ has form

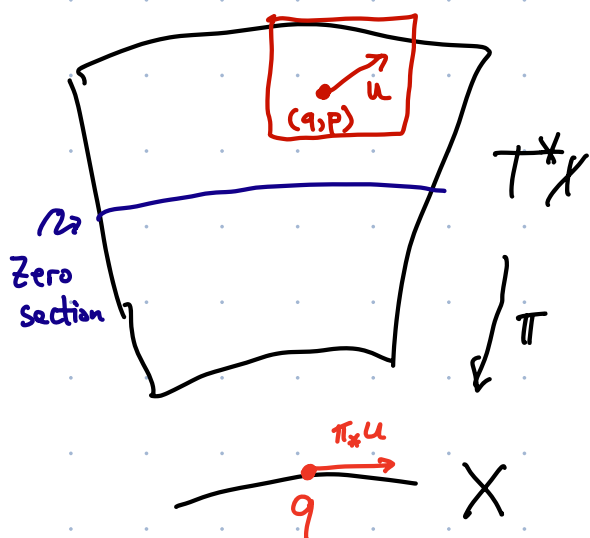
$$p_1 dx^1 + \dots + p_n dx^n$$

Phase space.

$\Rightarrow (x^1, \dots, x^n, p_1, \dots, p_n)$ coordinates on $T^*X = M$

Geometry of phase space:

The canonical 1-form $\Theta \in \Omega^1(M)$



given a vector $u \in T_{(q,p)} M$

$$\Theta(u) = p(\pi_* u)$$

\nwarrow derivative of π
 \nwarrow vector at q
 \nwarrow covector at q

in coords:

$$u = u^i \frac{\partial}{\partial q^i} + u_i \frac{\partial}{\partial p_i}$$

$$p = p_i dq^i$$

$$\Theta = \underbrace{a_i}_{p_i} dq^i + \underbrace{b^j}_0 dp_j$$

$$\begin{cases} \Theta(u) = p_i dq^i \left(\pi_* \left(u^j \frac{\partial}{\partial q^j} + u_j \frac{\partial}{\partial p_j} \right) \right) \\ = p_i u^i \end{cases}$$

$$\Theta = p_i dq^i$$

$$\pi(p, q) = q \quad \pi_* : \begin{matrix} \frac{\partial}{\partial p} \mapsto 0 \\ \frac{\partial}{\partial q} \mapsto \frac{\partial}{\partial q} \end{matrix}$$

since $\textcircled{4} \in \Omega^1(M)$ can differentiate,

obtain 2-form Symplectic form on M

$$\omega = d\mathbb{H} = d(p_i dq^i) = dp_i \wedge dq^i$$

$$= dp_1 \wedge dq^1 + \dots + dp_n \wedge dq^n$$

indep. of orig. choice of coords q^i
 i.e. $\tilde{q}^i = \tilde{q}^i(q^1, \dots, q^n)$ given coord. change
 \tilde{p}_i canonically conj. to \tilde{q}^i
 and in $(\tilde{p}_i, \tilde{q}^i)$ coords for M
 $\omega = d\tilde{p}_i \wedge d\tilde{q}^i$

view ω as a skew-symm. bilinear form

$$\text{on } T_{(p,q)} M$$

$$u_1, u_2 \in T_{(p,q)} M \quad \omega(u_1, u_2) = -\omega(u_2, u_1)$$

has matrix in basis $\frac{\partial}{\partial q}, \frac{\partial}{\partial p}$ given by

[illegible]

$$\beta = \begin{pmatrix} \frac{\partial}{\partial p_1} \\ \vdots \\ \frac{\partial}{\partial p_n} \\ \frac{\partial}{\partial q^1} \\ \vdots \\ \frac{\partial}{\partial q^n} \end{pmatrix} \quad \begin{matrix} \text{order} \\ \text{basis} \end{matrix} \quad \downarrow$$

key properties of ω :

- ① skew-symmetric $\omega \in \Omega^2(M)$
- ② closed $d\omega = 0$
- ③ ω nondegenerate

$$\omega : TM \xrightarrow{\cong} T^*M \quad \text{isomorphism.}$$

$$u \longmapsto i_u \omega = \omega(u, -)$$

$$\frac{\partial}{\partial p_i} \mapsto i \frac{\partial}{\partial p_i} (dp_j \wedge dq^j) = \left(i \frac{\partial}{\partial p_i} dp_j \right) dq^j + (-1) dp_j \wedge \left(i \frac{\partial}{\partial p_i} dq^j \right)$$

$$\frac{\partial}{\partial q_i} \longmapsto -dp_i$$

Main Consequence: we obtain a mechanism

by which $f \in C^\infty(M, \mathbb{R})$ is converted

to a vector field $X_f \in \mathcal{X}(M)$

Def: The Hamiltonian vector field of f is given by:

$$C^\infty(M, \mathbb{R}) \longrightarrow \mathcal{X}(M)$$

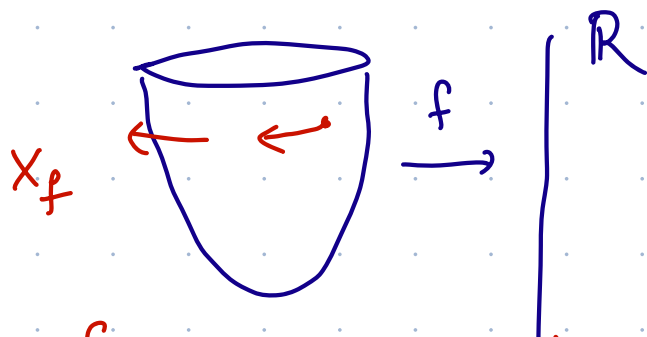
$$\begin{array}{ccc} f & \searrow d & \\ & df & \xrightarrow{-\omega^{-1}} -\omega^{-1}df = X_f \\ & \uparrow \pi & \\ & \Omega^1(M) & \end{array}$$

"Hamiltonian vector field associated to f ."

Main properties of Ham. v.f. are

Prop: The Ham. v.f. of f is st.

$$\textcircled{1} \quad X_f(f) = df(X_f) = df(-\omega^{-1}(df))$$



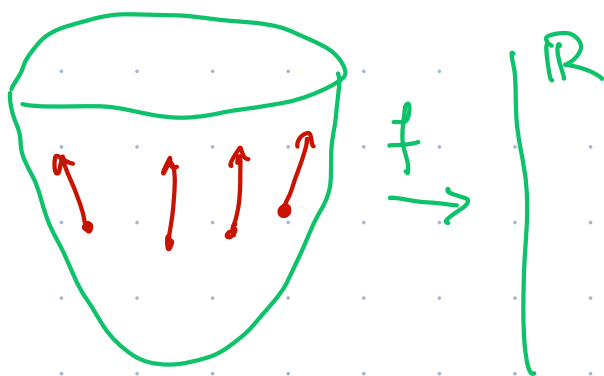
$$\begin{aligned} &= -\omega^{-1}(df, df) \\ &= 0 \quad \omega \text{ is skew} \end{aligned}$$

f is conserved by the flow of X_f

Warning: quite different from $\text{grad } f$

On a Riem. mfd $\nabla f = g^{-1}(df)$

$$\begin{aligned}(\nabla f)(f) &= df(g^{-1}df) = g^{-1}(df, df) \\ &= \|df\|^2\end{aligned}$$



$$\begin{aligned}\textcircled{2} \quad L_{X_f} \omega &= (di_{X_f} + i_{X_f} d) \omega \\ &\quad dw = 0 \\ &= d i_{X_f} \omega \\ &= d \left(\omega (-\omega^{-1} df, -) \right) \\ &= d(-df) = 0\end{aligned}$$

ω is preserved!

Prop: The hamiltonian v.f. of $\alpha f^\alpha f$
is a symmetry of (M, ω, f) .

Def: A Hamiltonian system (M, ω, H)
consists of

- ① A manifold M (for us $M = T^*X$)
- ② ω symplectic form on M
nondegenerate closed 2-form
- ③ A function H
"The Hamiltonian"

Theorem: Near any point p of a Symplectic
mfd, it is possible to find
"Darboux" coordinates $(q^1, \dots, q^n, p_1, \dots, p_n)$
s.t. $\omega = dp_i \wedge dq^i$

Note: There are degenerate versions of sym. forms

1. Poisson structure (ω^{-1} becomes degenerate)
2. quasi-symplectic (ω becomes deg.

↳ folded symplectic forms . . .

relevant to systems w/ constraints or singularities.

(Also there are higher degree variants
"Multi-symplectic" structures
 $\omega \in \Omega^{k \geq 2}(M)$)

Def: The Poisson bracket assoc. to ω is:

$$\{, \} : C^\infty(M, \mathbb{R}) \times C^\infty(M, \mathbb{R}) \longrightarrow C^\infty(M, \mathbb{R}).$$

$$\begin{aligned} \{f, g\} &= X_g(f) = df(-\omega^{-1}dg) \\ &= \omega^{-1}(df, dg). \\ &= -\omega^{-1}(dg, df) \\ &= -X_f(g) \end{aligned}$$

Prop: Poisson bracket satisfies Leibniz rule:

$$\begin{aligned}\{f_1 f_2, g\} &= X_g(f_1 f_2) = (X_g(f_1))f_2 + f_1(X_g(f_2)) \\ &= \{f_1, g\} f_2 + f_1 \{f_2, g\}\end{aligned}$$

Def: A Poisson algebra is a commutative \mathbb{R} -algebra A together w/ a bracket $\{, \} : A \times A \rightarrow A$

st. ① skew

② Leibniz $\{f, gh\} = \{f, g\}h + g\{f, h\}$

③ Jacobi identity

$$\{\{f, g\}, h\} = \{f, \{g, h\}\} - \{g, \{f, h\}\}.$$

Verify that our Poisson bracket satisfies Jacobi identity.

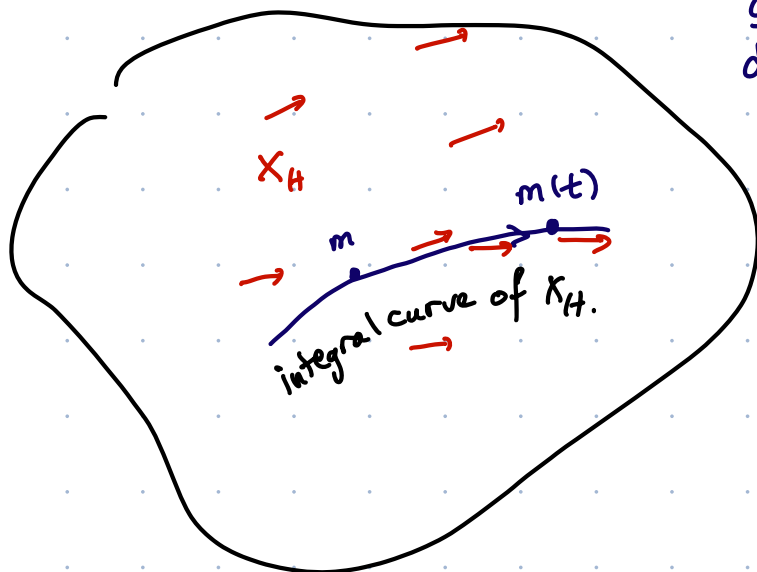
Poisson bracket is helpful for describing Hamiltonian flow

$$C^\infty(M) \ni H \xrightarrow{-\omega^\flat} X_H \quad \text{generates a flow}$$

$$\varphi_t^{X_H} : (\mathbb{R}, t) \times M \longrightarrow M$$

$$(t, m) \longmapsto \varphi_t^{X_H}(m) = m(t)$$

$$\frac{d}{dt} m(t) = X_H(m(t))$$



Alternatively we can flow a f^n $f \in C^\infty(M, \mathbb{R})$

i.e. $f(t) = (\varphi_t^{X_H})^* f$

$$\boxed{\frac{d}{dt} f(t) = X_H(f)} = \{f, H\}$$

Equation of motion for a f^n f dragged along

Ham. flow is

$$\frac{d}{dt} f = \{f, H\}$$

Hamilton's equation

Note: In any Poisson algebra $(A, \{, \})$ can write similar eqⁿ: Fix $H \in A$

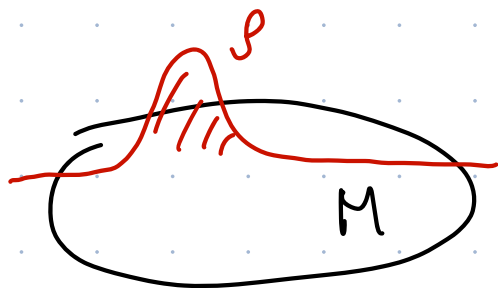
Then we can write a diff. eqn for a path

$f(t)$ of elts in A :

$$\dot{f} = \{f, H\}$$

If we know init. state m exactly \Rightarrow find evolution
 $\dot{m}(t) = X_H(m(t))$

If we only have probabilistic inform. about state



$$\dot{p} = \{p, H\}$$