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Introduction

Representations

Let $\mathfrak{g} = \mathfrak{sl}_n(\mathbb{C})$ be the Lie algebra of trace 0 matrices and $V = \mathbb{C}^n$. The standard basis vectors $\mathbf{v}_1, \ldots, \mathbf{v}_n$ form a basis for V that has several favorable properties:

Each basis vector is an eigenvector for the action of the subalgebra of diagonal matrices h, i.e. diag(t₁,...t_n) · **v**_k = t_k**v**_k

2 The matrices
$$E_{ij} = (e_{mn})$$
 s.t. $e_{mn} = \begin{cases} 1 & \text{if } (m, n) = (i, j) \\ 0 & \text{else} \end{cases}$

for $i \neq j$ "almost permute" these vectors, i.e. $E_{ij} \cdot \mathbf{v_j} = \mathbf{v_i}$ and $E_{ij} \cdot \mathbf{v_k} = \mathbf{0}$ for $k \neq j$.

In fact we only need to use the matrices $F_i = E_{i+1 \ i}$ to reach any basis vector from **v**₁.

We say that this is a **good basis**.

Introduction

Representations

Thus we can encode the representation as a colored directed graph, for example, \mathfrak{sl}_3 acting on \mathbb{C}^3 could be represented like this:

 $\mathbf{v_1} \stackrel{F_1}{\rightarrow} \mathbf{v_2} \stackrel{F_2}{\rightarrow} \mathbf{v_3}$

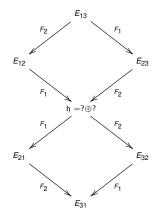
Our aim is to generalize this idea and we'd hope that the good basis is compatible with tensor product decompositions and branching.

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Introduction

Representations

This works splendidly only as long as each weight space is one-dimensional. We already run into trouble with the adjoint representation of \mathfrak{sl}_3 , as ker F_1 , ker F_2 , im F_1 , im F_2 are all different subspaces of \mathfrak{h} .



 Introduction

Quantized enveloping algebras

Consider the quantized enveloping algebra

$$\mathfrak{g} \rightsquigarrow U_q(\mathfrak{g})$$

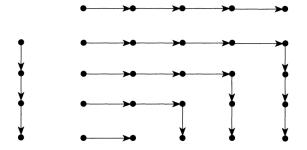
 $V \rightsquigarrow V_q$

- Solution 8 Sector 2 Kashiwara later developed crystals by going to the $q \rightarrow 0$ limit in the general case.

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Ordinary crystals

What is the benefit of crystals? Combinatorics. For $\mathfrak{g} = \mathfrak{sl}_2$ -crystals, tensor product decompositions are given by:



Ordinary crystals

We know that for an irreducible \mathfrak{sl}_n -representation V_λ of highest weight λ , dim $(V_\lambda) = \#SSYT(\lambda)$ with entries up to *n*. The crystal of V_λ has a realization as tableaux, for example, here is the crystal of the adjoint representation of \mathfrak{sl}_3

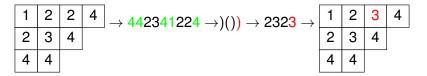


Ordinary crystals

The lowering operators f_i work as follows:

- Read the entries of the tableau from bottom to top, left to right, ignoring all numbers except i and i + 1.
- 2 Replace $i + 1 \rightarrow (\text{ and } i \rightarrow)$,
- Solution Turn the rightmost unmatched (into an i + 1

Example: applying f2



KR-crystals

Now consider the affine Lie algebra $\tilde{\mathfrak{g}}$. This is a central extension of the loop algebra $\mathfrak{g} \otimes \mathbb{C}[t, t^{-1}]$, we may consider $U(\tilde{\mathfrak{g}})$ and $U_q(\tilde{\mathfrak{g}})$.

Conjecture 1 (Hatayama et al. [4])

Certain finite-dimensional $U(\tilde{\mathfrak{g}})$ -modules, called **Kirillov-Reshetikhin modules** would have crystal bases. A KR-module $W_s^{(r)}$ is determined by a choice of a non-affine node (r) of the Dynkin diagram and a positive integer s.

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sl_n KR-crystals

As before, the situation is simplest for \mathfrak{sl}_n . It turns out that in this case, each KR-module stays irreducible when restricted to \mathfrak{sl}_n , for example, for $U_q(\mathfrak{sl}_3)$, W_1^1 is the standard representation, and the KR crystal is



Crystals	
sl_n KR-crystals	

Question: can we extend the tableau model to the KR setting? We want to give a combinatorial description of the f_0 operator. Shimozono's [6] method for \mathfrak{sl}_n : use Schützenberger's **promotion** operator on tableaux.

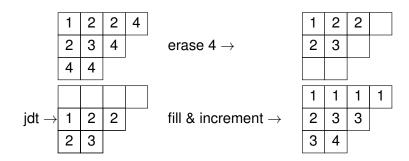
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If T is an SSYT, then pr(T) is obtained by the following procedure:

- Erase all entries *n* from the tableau.
- Icu-de-taquin the other boxes to the Southeast.
- Fill the empty boxes in the Northwest with zeros.
- Add one to each entry.

Crystals

sl_n KR-crystals



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Kirillov-Reshetikhin Crystals and Cacti	
Crystals	
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The promotion operator shifts the content of the tableau cyclically, and it is in fact cyclic of order *n* for rectangular tableaux, but not in general. Notice how *pr* interacts with the lowering operators for rectangular tableaux:

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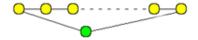
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sl_n KR-crystals

So the promotion operator realizes the cyclic symmetry of the affine type *A* Dynkin diagram, at least for rectangular tableaux.



Theorem 2 (Shimozono)

$$\mathit{f}_0 = \mathit{pr}^{-1} \circ \mathit{f}_1 \circ \mathit{pr}$$

One would think that this is a phenomenon specific to type *A* because of the cyclic symmetry, but there exist tableaux models for other types, where the KR crystals were shown to exist on a case-by-case basis by finding a suitable analog of the promotion operator [3].

The cactus group

The **cactus group** J_g was introduced by Henriques and Kamnitzer [8] in the context of coboundary categories. It has

- Generators: s_J for J a connected subdiagram of \mathfrak{g} 's Dynkin diagram.
- 2 Relations:

$$s_J^2 = 1 \ \forall J.$$

2
$$s_J s_{J'} = s_{J'} s_J$$
 if $J \cup J'$ is not connected.

where θ_J is the Dynkin diagram automorphism $-w_0^J$ of J. It surjects onto $W_{\mathfrak{g}}$ by $s_J \mapsto w_0^J$.

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The cactus group

Halacheva [7] showed that there is an action of J_g on any g-crystal, which we now describe.

The **Schützenberger involution** on a crystal B_{λ} of an irrep V_{λ} is the unique map $\xi_{\lambda} : B_{\lambda} \to B_{\lambda}$ on the vertices satisfying

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$$e_i(\xi_\lambda(b)) = \xi_\lambda(f_{\theta(i)}(b))$$

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$$f_i(\xi_\lambda(b)) = \xi_\lambda(e_{\theta(i)}(b))$$

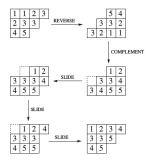
$$wt(\xi_{\lambda}(b)) = w_0 \cdot (wt(b))$$

So, in effect, it flips the crystal upside down.

The cactus group

For \mathfrak{sl}_n -crystals, the operation is given by **evacuation** on tableaux, which is the following procedure:

- Rotate the tableau 180°.
- 2 Complement the entries $i \rightarrow n + 1 i$.
- Jeu-de-taquin the tableau back to the original shape.

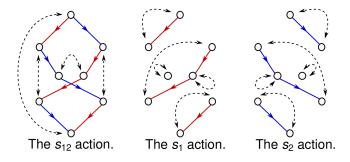


The cactus group

Halacheva [7] showed that J_g acts o a g-crystal B by

 $s_J(b) = \xi_{B_J}(b)$

where ξ_{B_J} is the Schützenberger involution on the restricted crystal using only the lowering operators in *J*. For example, on the adjoint representation of \mathfrak{sl}_3 ,



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Main result

Main theorem

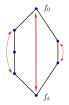
Our main result is the following:

Theorem 3

For a KR-module corresponding to a cominuscule fundamental weight (k), let J be the subset of g's Dynkin diagram complementary to $\{k\}$. Then

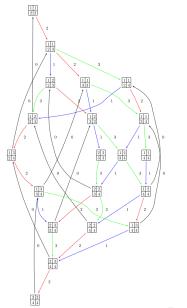
$$f_0^{-1}=e_0=s_Jf_ks_J.$$

In fact, s_J corresponds to



Main result

An example



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Main result

About the proof

Stembridge [9]:

- Check that $s_J e_k s_J$ is a lowering operator.
- See how $s_J e_k s_J$ interacts with the other lowering operators (e.g. $f_1(s_J e_k s_J)^2 f_1(b) = (s_J e_k s_J) f_1^2(s_J e_k s_J)(b)$ under certain local conditions).

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By the uniqueness of crystals, we are done.

Kirillov-Reshetikhin Crystals and Cacti	
Main result	
Remarks	

- Even though our method may be less ad hoc than trying to extend *pr* to other types, it only works for $W_s^{(r)}$ for (r) cominuscule.
- We are only using a single element of the cactus group, what do others do?
- When are other representations of g also representations of some Kac-Moody algebra containing g?
- What does this say about promotion and evacuation on Young tableaux?

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