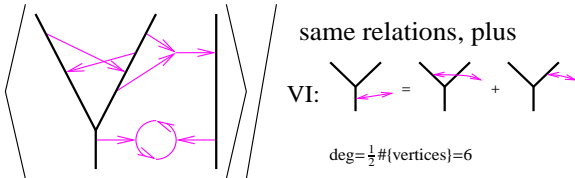
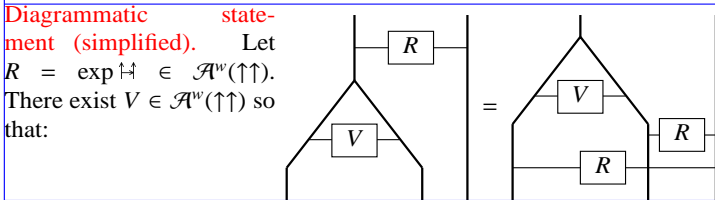


w-Jacobi diagrams and \mathcal{A} . $\mathcal{A}^w(Y \uparrow) \cong \mathcal{A}^w(\uparrow \uparrow \uparrow)$ is



Knot-Theoretic statement (simplified). There exists a homomorphic expansion Z for trivalent w-tangles. In particular, Z should respect R4.



Algebraic statement (simplified). With $r \in \mathfrak{g}^* \otimes \mathfrak{g}$ the identity element and with $R = e^r \in \hat{\mathcal{U}}(\mathfrak{g}) \otimes \hat{\mathcal{U}}(\mathfrak{g})$ there exist $V \in \hat{\mathcal{U}}(\mathfrak{g})^{\otimes 2}$ so that $V(\Delta \otimes 1)(R) = R^{13}R^{23}V$ in $\hat{\mathcal{U}}(\mathfrak{g})^{\otimes 2} \otimes \hat{\mathcal{U}}(\mathfrak{g})$

Unitary statement (simplified). There exists a unitary tangential differential operator V defined on $\text{Fun}(\mathfrak{g}_x \times \mathfrak{g}_y)$ so that $V e^{x+y} = e^x e^y V$ (allowing $\hat{\mathcal{U}}(\mathfrak{g})$ -valued functions)

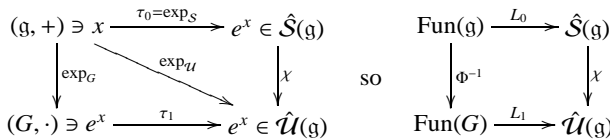
Group-Algebra statement (simplified). For every $\phi, \psi \in \text{Fun}(\mathfrak{g})^G$ (with small support), the following holds in $\hat{\mathcal{U}}(\mathfrak{g})$:

$$\iint_{\mathfrak{g} \times \mathfrak{g}} \phi(x)\psi(y)e^{x+y} = \iint_{\mathfrak{g} \times \mathfrak{g}} \phi(x)\psi(y)e^x e^y \quad (\text{shhh, this is Duflo})$$

Unitary \implies Group-Algebra. $\iint e^{x+y} \phi(x)\psi(y) = \langle 1, e^{x+y} \phi(x)\psi(y) \rangle = \langle V1, V e^{x+y} \phi(x)\psi(y) \rangle = \langle 1, e^x e^y V \phi(x)\psi(y) \rangle = \langle 1, e^x e^y \phi(x)\psi(y) \rangle = \iint e^x e^y \phi(x)\psi(y)$.

Convolutions statement (Kashiwara-Vergne, simplified). Convolutions of invariant functions on a Lie group agree with convolutions of invariant functions on its Lie algebra. More accurately, let G be a finite dimensional Lie group and let \mathfrak{g} be its Lie algebra, and let $\Phi : \text{Fun}(G) \rightarrow \text{Fun}(\mathfrak{g})$ be given by $\Phi(f)(x) := f(\exp x)$. Then if $f, g \in \text{Fun}(G)$ are Ad-invariant and supported near the identity, then $\Phi(f) \star \Phi(g) = \Phi(f \star g)$.

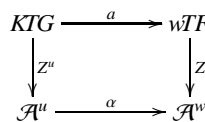
Convolutions and Group Algebras (ignoring all Jacobians). If G is finite, A is an algebra, $\tau : G \rightarrow A$ is multiplicative then $(\text{Fun}(G), \star) \rightarrow (A, \cdot)$ via $L : f \mapsto \sum f(a)\tau(a)$. For Lie (G, \mathfrak{g}) ,



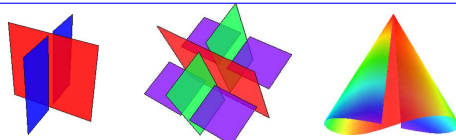
with $L_0 \psi = \int \psi(x) e^x dx \in \hat{\mathcal{S}}(\mathfrak{g})$ and $L_1 \Phi^{-1} \psi = \int \psi(x) e^x \in \hat{\mathcal{U}}(\mathfrak{g})$. Given $\psi_i \in \text{Fun}(\mathfrak{g})$ compare $\Phi^{-1}(\psi_1) \star \Phi^{-1}(\psi_2)$ and $\Phi^{-1}(\psi_1 \star \psi_2)$ in $\hat{\mathcal{U}}(\mathfrak{g})$:

$$\star \text{ in } G : \iint \psi_1(x)\psi_2(y)e^x e^y \quad \star \text{ in } \mathfrak{g} : \iint \psi_1(x)\psi_2(y)e^{x+y}$$

$u \leftrightarrow w$ The diagram on the right explains the relationship between associators and solutions of the Kashiwara-Vergne problem.



The Full
2-Knot Story



Question. Does it all extend to arbitrary 2-knots (not necessarily “simple”)? To arbitrary codimension-2 knots?

BF Following [CR]. $A \in \Omega^1(M = \mathbb{R}^4, \mathfrak{g})$, $B \in \Omega^2(M, \mathfrak{g}^*)$,

$$S(A, B) := \int_M \langle B, F_A \rangle.$$

With $\kappa : (S = \mathbb{R}^2) \rightarrow M$, $\beta \in \Omega^0(S, \mathfrak{g})$, $\alpha \in \Omega^1(S, \mathfrak{g}^*)$, set

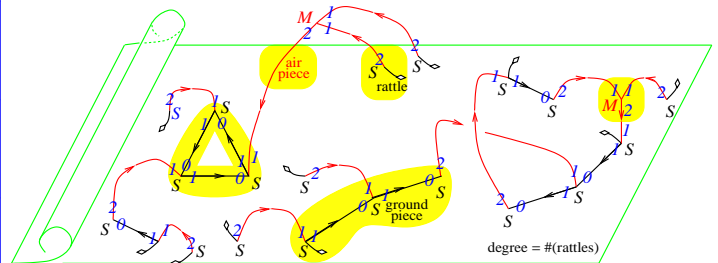
$$O(A, B, \kappa) := \int \mathcal{D}\beta \mathcal{D}\alpha \exp\left(\frac{i}{\hbar} \int_S \langle \beta, d_{\kappa^* A} \alpha + \kappa^* B \rangle\right).$$

The BF Feynman Rules. For an edge e , let Φ_e be its direction, in S^3 or S^1 . Let ω_3 and ω_1 be volume forms on S^3 and S^1 . Then

$$Z_{BF} = \sum_{\text{diagrams } D} \frac{[D]}{|\text{Aut}(D)|} \int_{\mathbb{R}^2} \dots \int_{\mathbb{R}^2} \int_{\mathbb{R}^4} \dots \int_{\mathbb{R}^4} \prod_{\text{red } e \in D} \Phi_e^* \omega_3 \prod_{\text{black } e \in D} \Phi_e^* \omega_1$$

(modulo some IHX-like relations).

See also [Wa]



Issues. • Signs don't quite work out, and BF seems to reproduce only “half” of the wheels invariant on simple 2-knots.

- There are many more configuration space integrals than BF Feynman diagrams and than just trees and wheels.
- I don't know how to define / analyze “finite type” for general 2-knots.
- I don't know how to reduce Z_{BF} to combinatorics / algebra.

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“God created the knots, all else in topology is the work of mortals.”

Leopold Kronecker (modified)

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