

LECTURES ON CURVE SHORTENING FLOW

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ABSTRACT. These are the lecture notes for the last three weeks of my PDE II course from Spring 2016. The curve shortening flow is a geometric heat equation for curves and provides an accessible setting to illustrate many important concepts from nonlinear PDE, including maximum principle estimates, monotonicity formulas, Harnack inequalities and blowup analysis. All these techniques will be combined to give an exposition of Huisken's proof of Grayson's beautiful theorem that the curve shortening flow shrinks any closed embedded curve in the plane to a round point.

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1. CURVE SHORTENING FLOW BASICS

A one-parameter family of embedded curves $\{\Gamma_t \subset \mathbb{R}^2\}_{t \in I}$ moves by curve shortening flow if the normal velocity at each point is given by the curvature vector:

$$(1.1) \quad \partial_t p = \vec{\kappa}(p)$$

for all $p \in \Gamma_t$ and all $t \in I$. Here, I is an interval, $\partial_t p$ is the normal velocity at p , and $\vec{\kappa}(p)$ is the curvature vector at p .

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Example 1.2 (Round shrinking circles). If $\Gamma_t = \partial B_{r(t)}^2 \subset \mathbb{R}^2$, then (1.1) reduces to an ODE for the radius, namely $\dot{r} = -1/r$. The solution with $r(0) = R$ is given by $r(t) = \sqrt{R^2 - 2t}$, $t \in (-\infty, R^2/2)$.

Example 1.3 (Grim reaper). Another explicit solution is given by $\Gamma_t = \text{graph}(-\log \cos p) + t$ where $p \in (-\frac{\pi}{2}, \frac{\pi}{2})$ and $t \in (-\infty, \infty)$.

Example 1.4 (Paperclip and hairclip). Despite the nonlinear nature of the curve shortening flow, the upwards translating grim reaper given by $e^{-y(t)} = e^{-t} \cos x(t)$ and the downwards translating grim reaper given by $e^{y(t)} = e^{-t} \cos x(t)$ can be combined to give another pair of solutions given implicitly as the solution set of

$$(1.5) \quad \cosh y(t) = e^{-t} \cos x(t),$$

respectively

$$(1.6) \quad \sinh y(t) = e^{-t} \cos x(t).$$

The paperclip, given as solution of (1.5) restricted to $|x| < \pi/2$ describes a compact ancient solution that for $t \rightarrow 0$ becomes extinct in a round point and for $t \rightarrow -\infty$ looks like two copies of the grim reaper glued together smoothly. The hairclip (1.6) is an eternal solution, which for $t \rightarrow -\infty$ looks like an infinite row of grim reapers, alternating between translating up and translating down, and for $t \rightarrow +\infty$ converges to a horizontal line.

From now on we will focus on the evolution of closed embedded curves. It is often most convenient to describe the evolution (1.1) in terms of a one-parameter family of embeddings

$$(1.7) \quad \gamma = \gamma(\cdot, t) : S^1 \times I \rightarrow \mathbb{R}^2$$

with $\Gamma_t = \gamma(S^1, t)$. Setting $p = \gamma(x, t)$, the curve shortening flow equation then takes the form

$$(1.8) \quad \partial_t \gamma(x, t) = \kappa(x, t) N(x, t),$$

where we have expressed the curvature vector $\vec{\kappa}$ as a product of the curvature κ and the inward pointing unit normal vector N .

Remark 1.9. The evolution can also be written in the form

$$(1.10) \quad \partial_t \gamma = \partial_s^2 \gamma,$$

where s denotes arc length. Even though this almost looks like the linear heat equation, the curve shortening flow is of course a nonlinear PDE since the arc length s depends in a nonlinear way on x and t .

We will for now assume that we have a smooth solution of (1.8) on a time interval $[0, T)$ and derive various elementary consequences.¹

Let $A(t)$ be the area of the domain enclosed by Γ_t . Then

$$(1.11) \quad \frac{d}{dt}A(t) = - \int_{\Gamma_t} \kappa ds = -2\pi,$$

and thus $A(t) = A(0) - 2\pi t$. In particular, $T \leq A(0)/2\pi$. We will see at the end of the semester (and this is a big theorem due to Grayson [Gra87], combined with earlier work in the convex case by Gage-Hamilton [GH86]) that the evolution actually can be continued smoothly until time $A(0)/2\pi$, where the flow becomes extinct in a round point.

A geometric incarnation of the maximum principle is the following comparison principle: If $\{\Gamma_t\}_{t \in [t_0, t_1]}$ and $\{\Gamma'_t\}_{t \in [t_0, t_1]}$ are two curve shortening flows (say at least one of them compact) which are initially disjoint, then they stay disjoint under the evolution, i.e.

$$(1.12) \quad \Gamma_{t_0} \cap \Gamma'_{t_0} = \emptyset \quad \Rightarrow \quad \Gamma_t \cap \Gamma'_t = \emptyset \quad \forall t \geq t_0.$$

In fact, it is easy to check that $\text{dist}(\Gamma_t, \Gamma'_t)$ is nondecreasing in time. In particular, any closed embedded curve which is contained in B_1 , even a curve spiraling around millions of times, will unwind itself and become round in time $T \leq \frac{1}{2}$. This vividly illustrates the strength of Grayson's theorem and the power of geometric heat equations in general.

We will next derive the evolution equation for the length $L(t) = \int_{\Gamma_t} ds$. To this end, we start by computing

$$(1.13) \quad \frac{d}{dt}L(t) = \frac{d}{dt} \int_{S^1} \langle \partial_x \gamma, \partial_x \gamma \rangle^{1/2} dx$$

$$(1.14) \quad = \int_{S^1} \langle \partial_x \partial_t \gamma, T \rangle dx$$

$$(1.15) \quad = \int_{S^1} \langle \partial_x (\kappa N), T \rangle dx,$$

where $T = \partial_x \gamma / |\partial_x \gamma|$ denotes the unit tangent and where we used (1.8) in the last step. Since $\langle N, T \rangle = 0$ and since $\langle \partial_x N, T \rangle = -\kappa \frac{ds}{dx}$ by the

¹The assumption that the initial curve is smooth is not really necessary. As a consequence of the estimates we'll discuss any C^2 -curve becomes instantaneously smooth under the curve shortening flow. In fact, the Cauchy problem for the curve shortening flow is well posed starting with any finite length Jordan curve [Lau13].

definition of curvature, we conclude that

$$(1.16) \quad \frac{d}{dt}L(t) = - \int_{\Gamma_t} \kappa^2 ds.$$

Equation (1.16) has the interpretation that the curve shortening flow is the gradient flow of the length functional, i.e. indeed deserves its name since it shortens curves as efficiently as possible (in general the variation of the length functional with normal speed v is given by $-\int_{\Gamma} \kappa v ds$).

Proposition 1.17 (Evolution equation for the curvature). *If $\{\Gamma_t \subset \mathbb{R}^2\}$ evolves by curve shortening flow, then its curvature evolves by*

$$(1.18) \quad \kappa_t = \kappa_{ss} + \kappa^3,$$

where s denotes arc length.

Proof. For convenience we work with a parametrization that satisfies $|\partial_x \gamma| = 1$ and $\langle \partial_x^2 \gamma, T \rangle = 0$ at the point (x, t) under consideration. Since $\kappa = |\partial_x \gamma|^{-2} \langle \partial_x^2 \gamma, N \rangle$, at the point (x, t) we can compute

$$(1.19) \quad \kappa_t = \partial_t \langle \partial_x^2 \gamma, N \rangle - 2 \langle T, \partial_x \partial_t \gamma \rangle \langle \partial_x^2 \gamma, N \rangle$$

$$(1.20) \quad = \langle \partial_x^2 \partial_t \gamma, N \rangle - 2\kappa \langle T, \partial_x \partial_t \gamma \rangle,$$

since $\partial_t N$ is tangential and $\partial_x^2 \gamma$ is normal. Using equation (1.8) for $\partial_t \gamma$ and using $\langle \partial_x N, N \rangle = 0$ we can continue our computation at (x, t) and obtain

$$(1.21) \quad \partial_t \kappa = \partial_x^2 \kappa + \kappa \langle \partial_x^2 N, N \rangle - 2\kappa^2 \langle T, \partial_x N \rangle$$

$$(1.22) \quad = \partial_x^2 \kappa - \kappa \langle \partial_x N, \partial_x N \rangle + 2\kappa^3.$$

Noting that $\partial_x^2 \kappa = \kappa_{ss}$ and $\langle \partial_x N, \partial_x N \rangle = \kappa^2$ at the point (x, t) , this proves the proposition. \square

Using Proposition 1.17 and the maximum principle we obtain:

Corollary 1.23. *Convexity is preserved under curve shortening flow, i.e. if $\kappa > 0$ at $t = 0$ then $\kappa > 0$ for all $t \in [0, T)$.*

More precisely, if $\kappa_{\min}(t) := \min_{\Gamma_t} \kappa$ is positive at $t = 0$, then it is nondecreasing in time and satisfies

$$(1.24) \quad \kappa_{\min}(t) \geq \frac{\kappa_{\min}(0)}{1 - 2t\kappa_{\min}^2(0)}.$$

In particular, this gives the (non sharp) estimate $T \leq 1/(2\kappa_{\min}^2(0))$.

We finish this first lecture by proving that the supremum of the curvature controls all the higher derivatives:

Theorem 1.25 (Derivative estimates). *There exist constants $C_\ell = C_\ell(K, T) < \infty$ such that if $\{\Gamma_t\}$ is a solution of the curve shortening flow with $\sup_{t \in [0, T)} \sup_{\Gamma_t} |\kappa| \leq K$, then*

$$(1.26) \quad \sup_{\Gamma_t} |\partial_s^\ell \kappa| \leq \frac{C_\ell}{t^{\ell/2}}.$$

Proof (sketch). Using the evolution equation (1.18) we compute

$$(1.27) \quad (\partial_t - \partial_s^2) \kappa^2 = -2\kappa_s^2 + 2\kappa(\partial_t - \partial_s^2) \kappa = -2\kappa_s^2 + 2\kappa^4.$$

Next, differentiating (1.18) with respect to arc length we obtain

$$(1.28) \quad (\kappa_t)_s = \kappa_{sss} + 3\kappa^2 \kappa_s,$$

which together with the commutator identity $(\kappa_s)_t = (\kappa_t)_s + \kappa^2 \kappa_s$ implies

$$(1.29) \quad (\partial_t - \partial_s^2) \kappa_s = 4\kappa^2 \kappa_s,$$

and thus

$$(1.30) \quad (\partial_t - \partial_s^2) \kappa_s^2 = -2\kappa_{ss}^2 + 8\kappa^2 \kappa_s^2.$$

Combining the evolution equations (1.27) and (1.30) we obtain

$$(1.31) \quad (\partial_t - \partial_s^2)(t\kappa_s^2 + \beta\kappa^2) \leq (8t\kappa^2 + 1 - 2\beta)\kappa_s^2 + 2\beta\kappa^4 \leq 2\beta K^4,$$

provided we chose $\beta \geq (8TK^2 + 1)/2$. Thus, the maximum principle implies that

$$(1.32) \quad t\kappa_s^2 \leq \beta K^2 + 2\beta K^4 T,$$

which proves the derivative estimate for $\ell = 1$.

For general ℓ , by induction one obtains the differential inequality

$$(1.33) \quad (\partial_t - \partial_s^2) |\partial_s^\ell \kappa|^2 \leq -2|\partial_s^{\ell+1} \kappa|^2 + \alpha_\ell \left(\sum_{i+j+k=\ell} |\partial_s^i \kappa| |\partial_s^j \kappa| |\partial_s^k \kappa| \right) |\partial_s^\ell \kappa|,$$

where α_ℓ are some numerical constants. The derivative estimates then follow by considering the evolution of

$$(1.34) \quad F_\ell = t^\ell |\partial_s^\ell \kappa|^2 + \sum_{i=0}^{\ell-1} \beta_{\ell,i} t^i |\partial_s^i \kappa|^2$$

for suitable constants $\beta_{\ell,i}$ and arguing by induction. The details are left as an exercise. \square

2. EXISTENCE AND UNIQUENESS

The goal of this lecture is to explain how to prove existence and uniqueness for the curve shortening flow.

Theorem 2.1 (Existence and Uniqueness). *Let $\gamma_0 : S^1 \rightarrow \mathbb{R}^2$ be an embedded curve. Then there exists a unique smooth solution $\gamma : S^1 \times [0, T) \rightarrow \mathbb{R}^2$ of the curve shortening flow*

$$(2.2) \quad \partial_t \gamma = \partial_s^2 \gamma, \quad \gamma|_{t=0} = \gamma_0,$$

defined on a maximal time interval $[0, T)$. The maximal existence time is characterized by

$$(2.3) \quad \sup_{S^1 \times [0, T)} |\kappa|(x, t) = \infty.$$

We start by explaining that the curve shortening flow is not strictly parabolic. If $\gamma : S^1 \times [0, T) \rightarrow \mathbb{R}^2$ evolves by curve shortening flow then

$$(2.4) \quad \partial_t \gamma = \partial_s^2 \gamma = |\partial_x \gamma|^{-1} \partial_x (|\partial_x \gamma|^{-1} \partial_x \gamma)$$

$$(2.5) \quad = |\partial_x \gamma|^{-2} \left(\partial_x^2 \gamma - \left\langle \frac{\partial_x \gamma}{|\partial_x \gamma|}, \partial_x^2 \gamma \right\rangle \frac{\partial_x \gamma}{|\partial_x \gamma|} \right).$$

In components $\gamma = (\gamma^1, \gamma^2)$ this reads

$$\begin{pmatrix} \partial_t \gamma^1 \\ \partial_t \gamma^2 \end{pmatrix} = |\partial_x \gamma|^{-4} \begin{pmatrix} (\partial_x \gamma^2)^2 & \partial_x \gamma^1 \partial_x \gamma^2 \\ \partial_x \gamma^1 \partial_x \gamma^2 & (\partial_x \gamma^1)^2 \end{pmatrix} \begin{pmatrix} \partial_x^2 \gamma^1 \\ \partial_x^2 \gamma^2 \end{pmatrix}.$$

Note that the matrix is positive semidefinite, but not positive definite (it has vanishing determinant). Thus the standard theory for strictly parabolic systems cannot be applied.

Degeneracies like above are actually quite typical for geometric PDEs. For the curve shortening flow, the underlying reason is that geometrically only the normal component of the velocity is meaningful, i.e. the velocity is only determined up to tangential motion / change of parametrization. To overcome this degeneracy we have to fix a gauge. To this end, we represent the flow as evolving graph over γ_0 (if γ_0 is not smooth, one can instead choose a smooth curve nearby), i.e. we make the ansatz

$$(2.6) \quad \tilde{\gamma}(x, t) = \gamma_0(x) + u(x, t)N(x).$$

Assume for convenience that γ_0 is parametrized by arc length. We compute

$$(2.7) \quad \tilde{\gamma}' = u'N + (1 - ku)T,$$

where k is the curvature of γ_0 , and T and N are the unit tangent and unit normal of γ_0 , respectively. Differentiating again we obtain

$$(2.8) \quad \tilde{\gamma}'' = (u'' + k(1 - ku))N - (k'u + 2ku')T.$$

Using the formula $\kappa = |\tilde{\gamma}'|^{-3} \det(\tilde{\gamma}', \tilde{\gamma}'')$ we obtain the curvature of $\tilde{\gamma}$,

$$(2.9) \quad \kappa = \frac{(1 - ku)u'' + 2ku'^2 + k'uu' - 2k^2u + k^3u^2 + k}{((1 - ku)^2 + u'^2)^{3/2}}.$$

The unit normal vector \tilde{N} of $\tilde{\gamma}$ is obtained by rotating $\tilde{\gamma}'/|\tilde{\gamma}'|$ by $\pi/2$ and thus equal to

$$(2.10) \quad \tilde{N} = \frac{(1 - ku)N - u'T}{((1 - ku)^2 + u'^2)^{1/2}}.$$

Finally, observing that $\Gamma_t = \tilde{\gamma}(S^1, t)$ moves by curve shortening flow if and only if $\langle \tilde{N}, u'N \rangle = \kappa$, we obtain the evolution equation

$$(2.11) \quad u_t = \frac{u'' + (1 - ku)^{-1}(2ku'^2 + k'uu' - 2k^2u + k^3u^2 + k)}{((1 - ku)^2 + u'^2)}.$$

Equation (2.11) is a quasilinear strictly parabolic equation (as long as say $|ku| \leq 1/2$), and thus has a unique solution on some time interval $[0, \varepsilon)$, c.f. the previous PDE lectures (this can be done e.g by using Picard iteration combined with the theory for the inhomogenous linear heat equation and energy estimates in H^j for j sufficiently large).

In general, $\tilde{\gamma}$ only solves the equation up to tangential motion, i.e.

$$(2.12) \quad \partial_t \tilde{\gamma} = \kappa \tilde{N} + f \partial_x \tilde{\gamma},$$

for some function $f = f(x, t)$. To get a parametrization $\gamma(x, t)$ that literally solves the equation $\partial_t \gamma = \partial_s^2 \gamma$ we let

$$(2.13) \quad \gamma(x, t) = \tilde{\gamma}(\varphi_t(x), t),$$

where $\varphi_t : S^1 \rightarrow S^1$ is the unique solution of the family of ODEs

$$(2.14) \quad \frac{d}{dt} \varphi_t(x) = -f(x, t) \frac{\partial_x \tilde{\gamma}(x, t)}{\partial_x \tilde{\gamma}(\varphi_t(x), t)}, \quad \varphi_0(x) = x.$$

We have thus proved short time existence and uniqueness for the curve shortening flow. Note also that embeddedness is preserved by the maximum principle. Now let $\{\Gamma_t\}_{t \in [0, T)}$ be a solution on a maximal time interval $[0, T)$. If

$$(2.15) \quad \limsup_{t \rightarrow T} \sup_{\Gamma_t} |\kappa| < \infty,$$

then by Theorem 1.25 all derivatives of the curvatures are bounded also up to time T . We can thus pass to a smooth limit Γ_T , and applying

the short time existence result we can continue the evolution until time $T + \varepsilon$; this contradicts the fact that T was maximal and finishes the proof of Theorem 2.1.

3. HUISKEN'S MONOTONICITY FORMULA AND APPLICATIONS

Recall that under curve shortening flow the length evolves by

$$(3.1) \quad \frac{d}{dt} \int_{\Gamma_t} ds = - \int_{\Gamma_t} \kappa^2 ds.$$

However, since $\text{Length}(\lambda\Gamma) = \lambda\text{Length}(\Gamma)$, this is not that useful when considering blowup sequences with $\lambda \rightarrow \infty$. A great advance was made by Huisken, who discovered a scale invariant monotone quantity. To describe this, let $\{\Gamma_t \subset \mathbb{R}^2\}$ be a curve shortening flow (say of closed curves, or complete curves with at most polynomial length growth), let $X_0 = (x_0, t_0)$ be a point in space-time, and let

$$(3.2) \quad \rho_{X_0}(x, t) = (4\pi(t_0 - t))^{-1/2} e^{-\frac{|x-x_0|^2}{4(t_0-t)}} \quad (t < t_0),$$

be the 1-dimensional backwards heat kernel centered at X_0 .

Theorem 3.3 (Huisken's monotonicity formula [Hui90]).

$$(3.4) \quad \frac{d}{dt} \int_{\Gamma_t} \rho_{X_0} ds = - \int_{\Gamma_t} \left| \kappa + \frac{\langle \gamma, N \rangle}{2(t_0 - t)} \right|^2 \rho_{X_0} ds \quad (t < t_0).$$

Huisken's monotonicity formula (3.4) can be thought of as weighted version of (3.1). A key property is its invariance under parabolic rescaling (cf. homework). Moreover, the equality case of (3.4) exactly characterizes the selfsimilarly shrinking solutions (cf. homework).

Proof. Without loss of generality $X_0 = (0, 0)$. The proof essentially amounts to deriving belows pointwise identity (3.7) for $\rho = \rho_0$.

Since the tangential gradient of ρ is given by $\partial_s \rho = \nabla \rho - \langle \nabla \rho, N \rangle N$, the intrinsic Laplacian of ρ can be expressed as

$$(3.5) \quad \partial_s^2 \rho = \langle T, \nabla_T \partial_s \rho \rangle = \langle T, \nabla_T \nabla \rho \rangle + \kappa \langle N, \nabla \rho \rangle.$$

Observing also that $\frac{d}{dt} \rho = \partial_t \rho + \kappa \langle N, \nabla \rho \rangle$, we compute

$$(3.6) \quad \begin{aligned} \left(\frac{d}{dt} + \partial_s^2 \right) \rho &= \partial_t \rho + \langle T, \nabla_T \nabla \rho \rangle + 2\kappa \langle N, \nabla \rho \rangle \\ &= \partial_t \rho + \langle T, \nabla_T \nabla \rho \rangle + \frac{\langle N, \nabla \rho \rangle^2}{\rho} - \left| \kappa - \frac{\langle N, \nabla \rho \rangle}{\rho} \right|^2 \rho + \kappa^2 \rho. \end{aligned}$$

We can now easily check that $\partial_t \rho + \langle T, \nabla_T \nabla \rho \rangle + \frac{\langle N, \nabla \rho \rangle^2}{\rho} = 0$. Thus

$$(3.7) \quad \left(\frac{d}{dt} + \partial_s^2 - \kappa^2 \right) \rho = - \left| \kappa - \frac{\langle \gamma, N \rangle}{2t} \right|^2 \rho.$$

Using also the evolution equation $\frac{d}{dt} ds = -\kappa^2 ds$, we conclude that

$$(3.8) \quad \frac{d}{dt} \int_{\Gamma_t} \rho ds = - \int_{\Gamma_t} \left| \kappa - \frac{\langle \gamma, N \rangle}{2t} \right|^2 \rho ds \quad (t < 0).$$

This proves the theorem. \square

We will now explain how Huisken's monotonicity formula can be used to study singularities via blowup analysis. Let $\{\Gamma_t \subset \mathbb{R}^2\}_{t \in [0, T]}$ be a curve shortening flow of closed embedded curves, defined on a maximal time interval $[0, T)$. From the previous lecture we know that the curvature blows up at the singular time, i.e.

$$(3.9) \quad \limsup_{t \rightarrow T} \max_{\Gamma_t} |\kappa| = \infty.$$

In the following we will assume that the singularity forms with the so called type I blowup rate

$$(3.10) \quad \max_{\Gamma_t} |\kappa| \leq \frac{C}{\sqrt{T-t}},$$

for some $C < \infty$ (this assumption will be justified in later lectures).

We say that $x_0 \in \mathbb{R}^2$ is a blowup point if there are sequences $t_i \rightarrow T$, $p_i \in \Gamma_{t_i}$ such that $|\kappa|(p_i) \rightarrow \infty$ and $p_i \rightarrow x_0$. By (3.9) there indeed exists a blowup point x_0 . We now rescale parabolically with center (x_0, T) , i.e. for $\lambda > 0$ consider the rescaled flow

$$(3.11) \quad \Gamma_t^\lambda := \lambda \cdot (\Gamma_{T+\lambda^{-2}t} - x_0), \quad t \in [-\lambda^2 T, 0).$$

Claim 3.12. *For $\lambda \rightarrow \infty$ the flows $\{\Gamma_t^\lambda\}_{t \in [-\lambda^2 T, 0)}$ converge smoothly to the family of round shrinking circles $\{\partial B_{\sqrt{-2t}}(0)\}_{t \in (-\infty, 0)}$.*

Proof. Rescaling the blowup rate (3.10) gives

$$(3.13) \quad \max_{\Gamma_t^\lambda} |\kappa| \leq \frac{C}{\sqrt{-t}}, \quad t \in [-\lambda^2 T, 0).$$

In particular, given any compact time interval $I \subset (-\infty, 0)$, the flow $\{\Gamma_t^\lambda\}$ is defined on I for λ large enough, and has uniformly bounded curvature on I . By the derivative estimates (Theorem 1.25) we also have locally uniform bounds for all the derivatives of the curvatures. Thus for any sequence $\lambda_i \rightarrow \infty$ we can find a subsequence λ_{i_k} such that $\{\Gamma_t^{\lambda_{i_k}}\}$ converges smoothly to a limit $\{\Gamma_t\}_{t \in (0, \infty)}$.

We will now analyze the limit $\{\Gamma_t\}_{t \in (0, \infty)}$: By construction, the limit is an ancient solution of the curve shortening flow. The limit is embedded with multiplicity one (this follows e.g. from the quantitative embeddedness estimate from Lecture 5). Using the definition of blowup point and comparison with round shrinking circles we see that $\Gamma_{-1}^\lambda \cap B_2(x_0) \neq \emptyset$ for λ large enough. Thus, the limit is nonempty. By Huisken's monotonicity formula for every $t_1 < t_2 < 0$ we have

$$(3.14) \quad \int_{t_1}^{t_2} \int_{\Gamma_t^\lambda} \left| \kappa - \frac{\langle \gamma, N \rangle}{2t} \right|^2 \rho ds dt = \left[- \int_{\Gamma_t} \rho_{(x_0, T)} ds \right]_{T-t_1/\lambda^2}^{T-t_2/\lambda^2} \rightarrow 0$$

as $\lambda \rightarrow \infty$. Thus, $\{\Gamma_t\}_{t \in (0, \infty)}$ is selfsimilarly shrinking and completely determined by its time slice Γ_{-1} which satisfies

$$(3.15) \quad \kappa + \frac{\langle \gamma, N \rangle}{2} = 0.$$

Since x_0 is a blowup point, Γ_{-1} cannot be a straight line by the local regularity theorem (see Theorem 3.17). Thus, Γ_{-1} must be a round circle of radius $\sqrt{2}$ (c.f. homework).

Finally, by uniqueness of the blowup limit, x_0 is unique and the subsequential convergence actually entails full convergence. \square

To finish this lecture, let us discuss the local regularity theorem which says if the density

$$(3.16) \quad \Theta(\{\Gamma_t\}, (x_0, t_0), r) := \int_{\Gamma_{t_0-r^2}} \rho_{(x_0, t_0)} ds$$

is close to one, then the curvature is controlled. We write $X = (x, t)$ for points in space-time and $P_r(X) = B_r(x) \times (t - r^2, t]$ for the parabolic ball with center X and radius r .

Theorem 3.17 (Local regularity theorem [Bra78, Whi05]). *There exist universal constants $\varepsilon > 0$ and $C < \infty$ with the following property: If $\{\Gamma_t \subset \mathbb{R}^2\}_{t \in (t_0-2r^2, t_0]}$ is a curve shortening flow (say of closed curves, or complete curves with at most polynomial length growth) with*

$$(3.18) \quad \sup_{\bar{X}_0 \in P_r(X_0)} \Theta(\{\Gamma_t\}, \bar{X}_0, r) < 1 + \varepsilon,$$

then

$$(3.19) \quad \sup_{P_{r/2}(X_0)} |\kappa| \leq Cr^{-1}.$$

Proof. Suppose the assertion fails. Then there exist a sequence of curve shortening flows $\{\Gamma_t^j \subset \mathbb{R}^2\}_{t \in (-2, 0]}$, with

$$(3.20) \quad \sup_{\bar{X}_0 \in P_1(0)} \Theta(\{\Gamma_t\}, \bar{X}_0, 1) < 1 + j^{-1},$$

but such that there are points $X_j \in P_{1/2}(0)$ with $|\kappa|(X_j) > j$.

By point selection, we can find $Y_j \in P_{3/4}(0)$ with $K_j = |\kappa|(Y_j) > j$ such that

$$(3.21) \quad \sup_{P_{j/(10K_j)}(Y_j)} |\kappa| \leq 2K_j.$$

Let us explain how the point selection works: Fix j . If $Y_j^0 = X_j$ already satisfies (3.21) with $K_j^0 = |\kappa|(Y_j^0)$, we are done. Otherwise, there is a point $Y_j^1 \in P_{j/(10K_j^0)}(Y_j^0)$ with $K_j^1 = |\kappa|(Y_j^1) > 2K_j^0$. If Y_j^1 satisfies (3.21), we are done. Otherwise, there is a point $Y_j^2 \in P_{j/(10K_j^1)}(Y_j^1)$ with $K_j^2 = |\kappa|(Y_j^2) > 2K_j^1$, etc. Note that $\frac{1}{2} + \frac{j}{10K_j^0}(1 + \frac{1}{2} + \frac{1}{4} + \dots) < \frac{3}{4}$. By smoothness, the iteration terminates after a finite number of steps, and the last point of the iteration lies in $P_{3/4}(0)$ and satisfies (3.21).

Continuing the proof of the theorem, let $\{\hat{\Gamma}_t^j\}$ be the flows obtained by shifting Y_j to the origin and parabolically rescaling by $K_j = |\kappa|(Y_j)$. Since the rescaled flow satisfies $|\kappa|(0) = 1$ and $\sup_{P_{j/10}(0)} |\kappa| \leq 2$, we can pass smoothly to a global limit. On the one hand, the limit is non-flat. On the other hand, by the rigidity case of Huisken's monotonicity formula and equation (3.20) the the limit is a straight line; a contradiction. \square

4. HAMILTON'S HARNACK INEQUALITY

Theorem 4.1 (Hamilton's Harnack inequality [Ham95]). *If $\{\Gamma_t \subset \mathbb{R}^2\}_{t \in [0, T]}$ is a convex solution of the curve shortening flow (say closed or complete with bounded curvature) then*

$$(4.2) \quad \frac{\kappa_t}{\kappa} - \frac{\kappa_s^2}{\kappa^2} + \frac{1}{2t} \geq 0.$$

Proof. The proof is very similar to the one of the Li-Yau Harnack. It suffices to consider the case where the solution existed since time $-\varepsilon$ and $\kappa \geq \varepsilon$. Let $f := \log \kappa$ and $F := t(f_s^2 - f_t)$. We want to use the maximum principle to show that $F \leq 1/2$ for all $t \in [0, T)$. Note that

$F \leq 1/2$ for small t . We compute

$$(4.3) \quad F_{ss} = t(2f_s f_{sss} + 2f_{ss}^2 - (f_t)_{ss})$$

$$(4.4) \quad = t(2f_s f_{sss} + 2f_{ss}^2 - (f_{ss})_t + 2\kappa^2 f_{ss} + 2\kappa^2 f_s^2),$$

where we used the commutator formula $[\partial_t, \partial_s] = \kappa^2 \partial_s$. Using the evolution equation for κ we see that $f_{ss} = -F/t - \kappa^2$ and thus

$$(4.5) \quad F_{ss} = -2f_s F_s + \frac{2F^2}{t} - \frac{F}{t} + F_t \\ - 4t\kappa^2 f_s^2 + 4F\kappa^2 + 2t\kappa^4 + 2t\kappa^2 f_t - 2\kappa^2 F - 2t\kappa^4 + 2t\kappa^2 f_s^2.$$

Miraculously, the nonlinear terms cancel and the quantity on the last line is identically zero, i.e.

$$(4.6) \quad F_{ss} - F_t = -2f_s F_s + \frac{1}{t} F(2F - 1)$$

If there is a maximum point (x_0, t_0) with $F(x_0, t_0) > 1/2$, then

$$(4.7) \quad 0 \geq (F_{ss} - F_t)|_{(x_0, t_0)} \geq 0 + \frac{1}{t} F(x_0, t_0)(2F(x_0, t_0) - 1) > 0;$$

a contradiction. This proves the theorem. \square

Applying the Harnack inequality with t replaced by $t - t_0$ and taking the limit $t_0 \rightarrow -\infty$ we obtain:

Corollary 4.8. *If $\{\Gamma_t \subset \mathbb{R}^2\}_{t \in (-\infty, T)}$ is an ancient convex solution of the curve shortening flow then*

$$(4.9) \quad \frac{\kappa_t}{\kappa} - \frac{\kappa_s^2}{\kappa^2} \geq 0,$$

in particular $\kappa_t \geq 0$.

Theorem 4.10 (Translating solitons [Ham95]). *Any eternal strictly convex solution $\{\Gamma_t \subset \mathbb{R}^2\}_{t \in (-\infty, \infty)}$ of the curve shortening flow such that κ has a critical point somewhere in space time, must be a translating soliton, i.e. there exists some vector $V \in \mathbb{R}^2$ such that $\Gamma_t = \Gamma_0 + tV$.*

Remark 4.11. The only strictly convex translating soliton, up to scaling and rigid motion, is the grim reaper (c.f. homework).

Proof. Assume κ has a critical point at (x_0, t_0) , i.e. $\kappa_t = 0 = \kappa_s$ at (x_0, t_0) . The Harnack quantity

$$(4.12) \quad Z = \frac{\kappa_t}{\kappa} - \frac{\kappa_s^2}{\kappa^2}$$

satisfies $Z \geq 0$ and $Z(x_0, t_0) = 0$. Using the strict maximum principle (c.f. the evolution equations in the above proof), we see that $Z \equiv 0$ for all $t \leq t_0$, i.e.

$$(4.13) \quad \kappa_t = \frac{\kappa_s^2}{\kappa}.$$

Consider the vector field

$$(4.14) \quad V := -\frac{\kappa_s}{\kappa}T + \kappa N.$$

We compute

$$(4.15) \quad V_s = \left(-\frac{\kappa_{ss}}{\kappa} + \frac{\kappa_s^2}{\kappa^2} - \kappa^2 \right) T + (\kappa_s - \kappa_s)N.$$

Using equation (4.13) we see that $V_s \equiv 0$. Similarly, $V_t \equiv 0$. Thus V is a constant vector. Since the normal component of V is given by κN , this implies that $\Gamma_t = \Gamma_{t_0} + (t - t_0)V$ for $t \leq t_0$. By uniqueness of the curve shortening flow we conclude that $\Gamma_t = \Gamma_0 + tV$ for all t . \square

5. HUISKEN'S DISTANCE COMPARISON PRINCIPLE

In this lecture, we discuss Huisken's estimate for the ratio between the intrinsic and extrinsic distance along the curve shortening flow. This can be thought of as quantitative version of the fact that embeddedness is preserved along the curve shortening flow. Let $X : S^1 \times [0, T) \rightarrow \mathbb{R}^2$ be a family of embedded curves evolving by curve shortening flow. Let $L(t)$ be the total length of the curve at time t . Given two points $x, y \in S^1$, denote by $\ell(x, y, t)$ the intrinsic distance between $X(x, t)$ and $X(y, t)$, and by $d(x, y, t) = |X(x, t) - X(y, t)|$ the extrinsic distance. Following Huisken, we consider the quantity

$$(5.1) \quad R(t) := \sup_{x \neq y} \frac{L(t)}{\pi d(x, y, t)} \sin \frac{\pi \ell(x, y, t)}{L(t)}.$$

Remark 5.2. Note that $R \geq 1$, and $R = 1$ only on the round circle.

Remark 5.3. Note that for $\ell \ll L$ we get $\frac{L}{\pi d} \sin \frac{\pi \ell}{L} \sim \frac{\ell}{d}$, i.e. the ratio between intrinsic and extrinsic distance. Since $\sin(\frac{\pi}{2} + \varphi) = \sin(\frac{\pi}{2} - \varphi)$, the function $\sin \frac{\pi \ell(x, y, t)}{L(t)}$ is smooth even at points with $\ell(x, y, t) = L(t)/2$.

Theorem 5.4 (Huisken's distance comparison principle [Hui98]). *If a family of closed embedded curves in the plane evolves by curve shortening flow, then the function $R(t)$ is nonincreasing in time.*

Proof. Let us describe the proof following [Hui98] and [Bre14]. If the assertion is false, we can find times $t_1 < t_2$ and a real number $r > 1$ such that

$$(5.5) \quad R(t_1) < r \quad \text{and} \quad R(t_2) > r.$$

We now consider the function

$$(5.6) \quad Z(x, y, t) := rd(x, y, t) - \frac{L(t)}{\pi} \sin \frac{\pi \ell(x, y, t)}{L(t)}.$$

By (5.5) there exists a time $\bar{t} \in (t_0, t_1)$ and a pair of points $\bar{x} \neq \bar{y}$ such that $Z(\bar{x}, \bar{y}, \bar{t}) = 0$, and $Z(x, y, t) \geq 0$ for all $x, y \in S^1$ and all $t \in (t_0, \bar{t})$.

Without loss of generality we can assume that the parametrization at time \bar{t} is by arc length, and that the orientation is chosen such that $\partial_x \ell(\bar{x}, \bar{y}, \bar{t}) = -1$ and $\partial_y \ell(\bar{x}, \bar{y}, \bar{t}) = +1$.

We start by computing the first derivatives

$$(5.7) \quad \frac{\partial Z}{\partial x}(x, y, \bar{t}) = r \frac{\langle X(x, \bar{t}) - X(y, \bar{t}), \frac{\partial X}{\partial x}(x, \bar{t}) \rangle}{|X(x, \bar{t}) - X(y, \bar{t})|} + \cos \frac{\pi \ell(x, y, \bar{t})}{L(\bar{t})},$$

and

$$(5.8) \quad \frac{\partial Z}{\partial y}(x, y, \bar{t}) = -r \frac{\langle X(x, \bar{t}) - X(y, \bar{t}), \frac{\partial X}{\partial y}(y, \bar{t}) \rangle}{|X(x, \bar{t}) - X(y, \bar{t})|} - \cos \frac{\pi \ell(x, y, \bar{t})}{L(\bar{t})}.$$

These first derivatives vanish when evaluated at $(\bar{x}, \bar{y}, \bar{t})$. In particular, adding these two identities gives

$$(5.9) \quad \left\langle X(\bar{x}, \bar{t}) - X(\bar{y}, \bar{t}), \frac{\partial X}{\partial x}(\bar{x}, \bar{t}) - \frac{\partial X}{\partial y}(\bar{y}, \bar{t}) \right\rangle = 0.$$

To keep the notation concise, we write $T(\bar{x}) = \frac{\partial X}{\partial x}(\bar{x}, \bar{t})$, $T(\bar{y}) = \frac{\partial X}{\partial y}(\bar{y}, \bar{t})$, $\kappa(\bar{x})N(\bar{x}) = \frac{\partial^2 X}{\partial x^2}(\bar{x}, \bar{t})$ and $\kappa(\bar{y})N(\bar{y}) = \frac{\partial^2 X}{\partial y^2}(\bar{y}, \bar{t})$, and also use the abbreviations $d = d(\bar{x}, \bar{y}, \bar{t})$, $\ell = \ell(\bar{x}, \bar{y}, \bar{t})$, $L = L(\bar{t})$, and

$$(5.10) \quad \omega = \frac{X(\bar{y}, \bar{t}) - X(\bar{x}, \bar{t})}{d}.$$

Using this notation, the identity (5.9) can be rewritten as

$$(5.11) \quad \langle \omega, T(\bar{x}) \rangle = \langle \omega, T(\bar{y}) \rangle.$$

We next compute the second order partial x -derivatives of Z :

$$\frac{\partial^2 Z}{\partial x^2}(\bar{x}, \bar{y}, \bar{t}) = \frac{r}{d} (1 - \langle \omega, T(\bar{x}) \rangle^2) - r\kappa(\bar{x}) \langle \omega, N(\bar{x}) \rangle + \frac{\pi}{L} \sin \frac{\pi \ell}{L}.$$

Similarly,

$$\frac{\partial^2 Z}{\partial y^2}(\bar{x}, \bar{y}, \bar{t}) = \frac{r}{d} (1 - \langle \omega, T(\bar{y}) \rangle^2) + r\kappa(\bar{y})\langle \omega, N(\bar{y}) \rangle + \frac{\pi}{L} \sin \frac{\pi\ell}{L},$$

and

$$\frac{\partial^2 Z}{\partial x \partial y}(\bar{x}, \bar{y}, \bar{t}) = -\frac{r}{d} (\langle T(\bar{x}), T(\bar{y}) \rangle - \langle T(\bar{x}), \omega \rangle \langle \omega, T(\bar{y}) \rangle) - \frac{\pi}{L} \sin \frac{\pi\ell}{L}.$$

Define $\alpha \in (0, \pi/2)$ by $\cos \alpha = \langle \omega, T(\bar{x}) \rangle = \langle \omega, T(\bar{y}) \rangle$. Then $\langle T(\bar{x}), T(\bar{y}) \rangle = \cos(2\alpha)$, and thus

$$(5.12) \quad \frac{\partial^2 Z}{\partial x^2}(\bar{x}, \bar{y}, \bar{t}) + \frac{\partial^2 Z}{\partial y^2}(\bar{x}, \bar{y}, \bar{t}) - 2 \frac{\partial^2 Z}{\partial x \partial y}(\bar{x}, \bar{y}, \bar{t}) = \\ - r\kappa(\bar{x})\langle \omega, N(\bar{x}) \rangle + r\kappa(\bar{y})\langle \omega, N(\bar{y}) \rangle + \frac{4\pi}{L} \sin \frac{\pi\ell}{L}.$$

Finally, the time derivative of Z can be computed as

$$\frac{\partial Z}{\partial t}(\bar{x}, \bar{y}, \bar{t}) = -r\langle \omega, \kappa(\bar{x})N(\bar{x}) - \kappa(\bar{y})N(\bar{y}) \rangle \\ + \left(\frac{1}{\pi} \sin \frac{\pi\ell}{L} - \frac{\ell}{L} \cos \frac{\pi\ell}{L} \right) \int_{S^1} \kappa^2 + \cos \frac{\pi\ell}{L} \int_{\bar{x}}^{\bar{y}} \kappa^2.$$

Since $r > 1$ and $Z(\bar{x}, \bar{y}, \bar{t}) = 0$, the curve $X(S^1, \bar{t})$ cannot have constant curvature, and thus

$$(5.13) \quad \int_{S^1} \kappa^2 > \frac{1}{L} \left(\int_{S^1} \kappa \right)^2 = \frac{4\pi^2}{L}.$$

Similarly,

$$(5.14) \quad \int_{\bar{x}}^{\bar{y}} \kappa^2 \geq \frac{1}{\ell} \left(\int_{\bar{x}}^{\bar{y}} \kappa \right)^2 = \frac{4\alpha^2}{\ell}.$$

Putting everything together, we conclude that

$$0 \geq \frac{\partial Z}{\partial t}(\bar{x}, \bar{y}, \bar{t}) - \frac{\partial^2 Z}{\partial x^2}(\bar{x}, \bar{y}, \bar{t}) - \frac{\partial^2 Z}{\partial y^2}(\bar{x}, \bar{y}, \bar{t}) + 2 \frac{\partial^2 Z}{\partial x \partial y}(\bar{x}, \bar{y}, \bar{t}) \\ > \frac{4}{\ell} \left(\alpha^2 - \frac{\pi^2 \ell^2}{L^2} \right) \cos \frac{\pi\ell}{L}.$$

On the other hand, the inequality $r > 1$ implies $\cos \alpha \leq \cos \frac{\pi\ell}{L}$, hence $\alpha \geq \frac{\pi\ell}{L}$. This is a contradiction. \square

Remark 5.15. By the monotonicity we have $R(t) \leq C$, where $C := R(0) < \infty$ measures the quantitative embeddedness of the initial curve. Note also that R is scaling invariant. In particular, this implies that the grim reaper and the paperclip cannot arise as a blowup limit of a curve shortening flow of closed embedded curves.

6. GRAYSON'S CONVERGENCE THEOREM

In this final lecture we explain Huisken's proof of Grayson's theorem:

Theorem 6.1 (Grayson's theorem [Gra87]). *If $\Gamma \subset \mathbb{R}^2$ is a closed embedded curve, then the curve shortening flow $\{\Gamma_t\}_{t \in [0, T]}$ with $\Gamma_0 = \Gamma$ exists until $T = \frac{A_\Gamma}{2\pi}$ and converges for $t \rightarrow T$ to a round point, i.e. there exists a unique point $x_0 \in \mathbb{R}^2$ such that the rescaled flows*

$$(6.2) \quad \Gamma_t^\lambda := \lambda \cdot (\Gamma_{T+\lambda^{-2}t} - x_0)$$

converge for $\lambda \rightarrow \infty$ to the round shrinking circle $\{\partial B_{\sqrt{-2t}}\}_{t \in (-\infty, 0)}$.

Remark 6.3. There are by now several different proofs of Grayson's theorem, in particular a nice geometric proof by Andrews-Bryan [AB11].

Lemma 6.4 (Grayson). *Along the curve shortening flow we have*

$$(6.5) \quad \frac{d}{dt} \int_{\Gamma_t} |\kappa| ds = -2 \sum_{x: \kappa(x, t) = 0} |\kappa_s|(x, t)$$

Proof of Lemma 6.4. Since solutions of the curve shortening flow are analytic, there are only finitely many inflection points. We compute

$$(6.6) \quad \frac{d}{dt} \left(\int_{\{\kappa \geq 0\}} \kappa ds - \int_{\{\kappa \leq 0\}} \kappa ds \right) = \int_{\{\kappa \geq 0\}} \kappa_{ss} ds - \int_{\{\kappa \leq 0\}} \kappa_{ss} ds.$$

Integrating by parts, the assertion follows. \square

Proof of Theorem 6.1. Let $T < \infty$ be the maximal existence time of the curve shortening flow starting at Γ . Suppose towards a contradiction that

$$(6.7) \quad \limsup_{t \rightarrow T} \left((T - t) \max_{\Gamma_t} \kappa^2 \right) = \infty.$$

We now perform a type II blowup as in [HS99], see also [Alt91]. For any integer $k \geq 1/T$ let $t_k \in [0, T - 1/k]$, $x_k \in S^1$ be such that

$$(6.8) \quad \kappa^2(x_k, t_k)(T - 1/k - t_k) = \max_{t \leq T - 1/k, x \in S^1} \kappa^2(x, t)(T - 1/k - t).$$

Furthermore we set

$$(6.9) \quad \lambda_k = \kappa(x_k, t_k), \quad t_k^{(0)} = -\lambda_k^2 t_k, \quad t_k^{(1)} = \lambda_k^2 (T - 1/k - t_k).$$

By (6.7), given any $M < \infty$ there exist $\bar{t} < T$ and $\bar{x} \in S^1$ such that $\kappa^2(\bar{x}, \bar{t})(T - \bar{t}) > 2M$. For k large enough we have

$$(6.10) \quad \bar{t} < T - 1/k, \quad \kappa^2(\bar{x}, \bar{t})(T - \bar{t} - 1/k) > M.$$

It follows that

$$(6.11) \quad t_k^{(1)} = \kappa^2(x_k, t_k)(T - 1/k - t_k) \geq \kappa^2(\bar{x}, \bar{t})(T - 1/k - \bar{t}) > M.$$

Since $t_k^{(1)}$ is increasing and M is arbitrary, this implies $t_k^{(1)} \rightarrow \infty$. It follows that $\lambda_k \rightarrow \infty$, $t_k \rightarrow T$ and $t_k^{(0)} \rightarrow -\infty$.

We now consider the sequence of rescaled flows

$$(6.12) \quad \Gamma_t^k = \lambda_k \cdot \left(\Gamma_{t_k + \lambda_k^{-2}t} - x_k \right), \quad t \in [t_k^{(0)}, t_k^{(1)}].$$

By construction, after choosing a suitable orientation, Γ_t^k has $\kappa_k = 1$ at $t = 0$ at the origin. In addition, our choice of (x_k, t_k) implies

$$(6.13) \quad \kappa_k^2(x, t) \leq \frac{T - 1/k - t_k}{T - 1/k - t_k - \lambda_k^2 t} = \frac{t_k^{(1)}}{t_k^{(1)} - t}, \quad t \in [t_k^{(0)}, t_k^{(1)}].$$

After passing to a subsequence, we thus get a smooth limit $\{\Gamma_t^\infty\}_{t \in (-\infty, \infty)}$. The limit satisfies $\kappa = 1$ at time 0 at the origin, and $\kappa^2 \leq 1$ at every point and every time. Moreover, by Lemma 6.4 the limit satisfies

$$(6.14) \quad \int_{-\infty}^{\infty} \sum_{x: \kappa(x, t) = 0} |\kappa_s|(x, t) dt = 0,$$

i.e. if there was any point with $\kappa = 0$ then we would have $\kappa_s = 0$ at this point also. Using the evolution equations and analyticity this would imply that $\{\Gamma_t^\infty\}_{t \in (-\infty, \infty)}$ is a straight line [Ang91]; a contradiction. Thus $\kappa > 0$, and by the equality case of Hamilton's Harnack inequality from Lecture 4 and the classification of translating solitons from the homework $\{\Gamma_t^\infty\}_{t \in (-\infty, \infty)}$ must be a grim reaper; this contradicts the bound for the ratio between the intrinsic and extrinsic distance from Lecture 5. We thus have proved the type I blowup rate bound

$$(6.15) \quad \limsup_{t \rightarrow T} \left((T - t) \max_{\Gamma_t} \kappa^2 \right) < \infty.$$

and by the results from Lecture 3 we conclude that $T = \frac{A\Gamma}{2\pi}$ and that for $t \rightarrow T$ the flow converges to a round point. \square

REFERENCES

- [AB11] B. Andrews and P. Bryan. Curvature bound for curve shortening flow via distance comparison and a direct proof of Grayson's theorem. *J. Reine Angew. Math.*, 653:179–187, 2011.
- [Alt91] S. Altschuler. Singularities of the curve shrinking flow for space curves. *J. Differential Geom.*, 34(2):491–514, 1991.
- [Ang91] S. Angenent. On the formation of singularities in the curve shortening flow. *J. Differential Geom.*, 33(3):601–633, 1991.

- [Bra78] K. Brakke. *The motion of a surface by its mean curvature*, volume 20 of *Mathematical Notes*. Princeton University Press, Princeton, N.J., 1978.
- [Bre14] S. Brendle. Two-point functions and their applications in geometry. *Bull. Amer. Math. Soc.*, 51(4):581–596, 2014.
- [GH86] M. Gage and R. S. Hamilton. The heat equation shrinking convex plane curves. *J. Differential Geom.*, 23(1):69–96, 1986.
- [Gra87] M. A. Grayson. The heat equation shrinks embedded plane curves to round points. *J. Differential Geom.*, 26(2):285–314, 1987.
- [Ham95] R. S. Hamilton. Harnack estimate for the mean curvature flow. *J. Differential Geom.*, 41(1):215–226, 1995.
- [HS99] G. Huisken and C. Sinestrari. Mean curvature flow singularities for mean convex surfaces. *Calc. Var. Partial Differential Equations*, 8(1):1–14, 1999.
- [Hui90] G. Huisken. Asymptotic behavior for singularities of the mean curvature flow. *J. Differential Geom.*, 31(1):285–299, 1990.
- [Hui98] G. Huisken. A distance comparison principle for evolving curves. *Asian J. Math.*, 2(1):127–133, 1998.
- [Lau13] J. Lauer. A new length estimate for curve shortening flow and low regularity initial data. *Geom. Funct. Anal.*, 23(6):1934–1961, 2013.
- [Whi05] B. White. A local regularity theorem for mean curvature flow. *Ann. of Math. (2)*, 161(3):1487–1519, 2005.