

Mat1062: Introductory Numerical Methods for PDE

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1 Ownership

These notes are the joint property of Rob Almgren and Mary Pugh.

2 Finite Element Methods

For finite element methods we need a Hilbert Space. This is a complete¹ vector space which has an inner product and a countable basis.

To start, we will focus on the two-point boundary value problem

$$\begin{cases} -u''(x) = f(x) & \forall x \in (0, 1) \\ u(0) = 0 \\ u'(1) = 0 \end{cases} \quad (1)$$

Given a continuous f , the solution is easy to find:

$$u(x) = F(x) + mx + b$$

where $F''(x) = f(x)$ and the integration constants m and b are chosen to satisfy the boundary conditions. We say that u is a *classical solution* if u has two derivatives, if u'' is continuous on $(0, 1)$ and if the boundary value

¹A space is complete if any convergent sequence converges to something *in* the space. For example, the real numbers are complete because any convergent sequence of real numbers converges to a real number. (They can't converge to a complex number, for example.) The rational numbers are not complete because there are convergent sequences of rational numbers that don't converge to a rational number. (One can construct a sequence of rational numbers that converges to $\sqrt{2}$, for example.)

problem is satisfied. We would like to define a weaker type of solution, one which is based on a variational formulation. (Because the problem is exactly solvable all of this may feel a little pointless. But the point is: we will be studying a simple case using methods that generalise to harder problems.)

To understand how we might formulate a weak solution, assume that u is a classical solution and v is a continuous function whose first derivative v' is continuous on $(0, 1)$. Then

$$\begin{aligned} \int_0^1 f(x)v(x) \, dx &= - \int_0^1 u''(x)v(x) \, dx \\ &= \int_0^1 u'(x)v'(x) \, dx - u'(1)v(1) + u'(0)v(0) \\ &= \int_0^1 u'(x)v'(x) \, dx + u'(0)v(0) \end{aligned}$$

Above, we used that $u'(1) = 0$. This calculation helps us in choosing what Hilbert Space we will look in for our solution. Specifically, we choose the Hilbert space

$$V = \{v(x) \text{ real-valued functions on } \mathbb{R} \mid \int_0^1 v(x)^2 \, dx < \infty, \int_0^1 v'(x)^2 \, dx < \infty, v(0) = 0\} \quad (2)$$

We then define a *weak solution* as: if $f \in L^2([0, 1])$ then u is a weak solution of the boundary value problem (1) if $u \in V$ and $\int u'v' = \int fv$ for all $v \in V$.

Before proceeding, we ask if the Hilbert space V makes sense. It's a subspace of $H^1([0, 1])$ which is a Hilbert space. And so it will be a Hilbert space in its own right. The only thing to worry about is whether it makes sense to specify the value of v at a point. If all we knew about a function, v , was that $\int v^2 < \infty$ then it would not make sense to ask what v equals at a particular point. However, it turns out that if $\int v'^2 < \infty$ then this forces v to be continuous on $[0, 1]$. And so it does make sense to talk about the value of v at a point. Note that the space V has one of the two boundary conditions built in to it. And so if $u \in V$ it automatically satisfies $u(0) = 0$. A natural question is whether or not we know that $u'(1) = 0$. (Note: if all we know is that $\int u'^2 < \infty$ then we can't know pointwise information like $u'(1) = 0$. But if we happen to know that u' is continuous then we would hope that $u'(1) = 0$ follows somehow.)

Now that we're satisfied with the definition of V we ask whether

$$\int_0^1 u'(x)v'(x) \, dx = \int_0^1 f(x)v(x) \, dx, \quad \forall v \in V$$

makes sense. Specifically, we need that each of these integrals is finite. This follows via the Schwartz inequality:

$$\left| \int_0^1 f(x)v(x) \, dx \right| \leq \sqrt{\int_0^1 f(x)^2 \, dx} \sqrt{\int_0^1 v(x)^2 \, dx} < \infty,$$

where we used that $f \in L^2([0, 1])$ and that $v \in V$. Similarly,

$$\left| \int_0^1 u'(x)v'(x) \, dx \right| \leq \sqrt{\int_0^1 u'(x)^2 \, dx} \sqrt{\int_0^1 v'(x)^2 \, dx} < \infty$$

where we used that $u, v \in V$.

Now that we're satisfied with our definition of weak solution, we ask the natural question: "If u is smooth and is a weak solution does this imply u is a classical solution?" The answer is "yes". (If the answer were "no" then we'd really have to question our definition of weak solution.) Before proving this, we introduce some notation:

$$a(u, v) := \int_0^1 u'(x)v'(x) \, dx \quad \langle f, v \rangle := \int_0^1 f(x)v(x) \, dx$$

Theorem *Assume f and u'' are continuous on $[0, 1]$. Let the space V be as defined in (2). If $u \in V$ and $a(u, v) = \langle f, v \rangle$ for all $v \in V$ then u is a classical solution of the boundary value problem (1).*

Proof Choose $v \in V$ such that v' is continuous on $[0, 1]$. Then

$$\begin{aligned} \int_0^1 u'(x)v'(x) \, dx &= \int_0^1 f(x)v(x) \, dx \\ \implies - \int_0^1 u''(x)v(x) \, dx + u'(1)v(1) - u'(0)v(0) &= \int_0^1 f(x)v(x) \, dx \\ \implies - \int_0^1 u''(x)v(x) \, dx + u'(1)v(1) &= \int_0^1 f(x)v(x) \, dx \\ \implies u'(1)v(1) &= \int_0^1 (f(x) + u''(x)) v(x) \, dx \end{aligned} \tag{3}$$

The identity (3) holds for any $v \in V$ that has a continuous first derivative. We now show that this implies that the continuous function $f + u''$ is identically zero on $[0, 1]$. Assume that $f + u''$ isn't identically zero on $[0, 1]$. Because $f + u''$ is continuous on $[0, 1]$ this would imply we could find an interval $(x_0, x_1) \subset [0, 1]$ on which $f + u''$ is positive. (If we can find no such interval then we can find an interval on which $f + u''$ is strictly negative.) We use this interval to construct a specific test function

$$v(x) = \begin{cases} (x - x_0)^2(x - x_1)^2 & x_0 \leq x \leq x_1 \\ 0 & \text{otherwise} \end{cases}$$

For this test function, (3) becomes

$$0 = \int_0^1 (f(x) + u''(x)) v(x) dx = \int_{x_0}^{x_1} (f(x) + u''(x)) v(x) dx > 0$$

which is impossible. (Note: if $f + u''$ were strictly negative on the interval then the above would have yielded $0 < 0$ which is again impossible.) This proves that $-u''(x) = f(x)$ at all points in $[0, 1]$. As a result, identity (3) reduces to

$$u'(1)v(1) = 0$$

for any $v \in V$ that has a continuous first derivative. For example, it holds for $v(x) = x$ resulting in $u'(1)v(1) = u'(1) = 0$. There was nothing really special about this choice — any $v \in V$ which has continuous first derivative (so that (3) applies) and has $v(1) \neq 0$ would have resulted in $u'(1) = 0$. This finishes the proof.

Essential boundary conditions and natural boundary conditions

We have boundary conditions at $x = 0$ and $x = 1$. The boundary condition at $x = 0$ is $u(0) = 0$. We built this boundary condition into the function space V . We're looking for a solution in the function space V — if we find a solution in V it will automatically satisfy the left boundary condition because that's part of the space's definition. When a boundary condition is part of how the function space is defined, the boundary condition is called *essential*. We just found that if a solution of the variational formulation is smooth enough (i.e. u'' is continuous) then it will satisfy $u'(1) = 0$. When a boundary condition follows from the variational formulation in this way, the boundary condition is called *natural*.

We now have a weak formulation that we are happy with:

$$\mathbf{u} \in \mathbf{V}, \quad \mathbf{a}(\mathbf{u}, \mathbf{v}) = \langle \mathbf{f}, \mathbf{v} \rangle \quad \forall \mathbf{v} \in \mathbf{V} \quad (4)$$

Another way to understand this would be

$$\mathbf{u} \in \mathbf{V}, \quad \int_0^1 \mathbf{u}'(x)\mathbf{v}'(x) - \mathbf{f}(x)\mathbf{v}(x) \, dx = 0 \quad \forall \mathbf{v} \in \mathbf{V}$$

which is the same thing as looking at the first variation in \mathbf{V} of the action functional

$$\mathcal{S}(\mathbf{u}) = \int_0^1 \frac{1}{2} \mathbf{u}_x^2(x) - \mathbf{f}(x)\mathbf{u}(x) \, dx$$

with Lagrangian $1/2 \mathbf{u}_x^2 - \mathbf{f}\mathbf{u}$. And so if you remember Hamilton's principle from classical mechanics you would see how the action functional would lead to the first variation would lead to the Euler-Lagrange equation which is $-\mathbf{u}_{xx} = \mathbf{f}$. If you're trying to solve a problem which has a Lagrangian then it is natural to use a finite element approach to implement Hamilton's principle using whatever family of basis functions you've chosen.

Finally, we note that $\mathbf{a}(\mathbf{u}, \mathbf{v})$ is actually an inner product on \mathbf{V} . To be an inner product we need to check the following

1. $\mathbf{a}(\mathbf{u}, \mathbf{v}) = \overline{\mathbf{a}(\mathbf{v}, \mathbf{u})} \quad \forall \mathbf{u}, \mathbf{v} \in \mathbf{V}$. This holds automatically because \mathbf{u} and \mathbf{v} are real-valued functions.
2. $\mathbf{a}(\mathbf{u}, \mathbf{v} + \mathbf{w}) = \mathbf{a}(\mathbf{u}, \mathbf{v}) + \mathbf{a}(\mathbf{u}, \mathbf{w}) \quad \forall \mathbf{u}, \mathbf{v}, \mathbf{w} \in \mathbf{V}$. This holds because $(\mathbf{v} + \mathbf{w})' = \mathbf{v}' + \mathbf{w}'$.
3. $\mathbf{a}(\lambda \mathbf{u}, \mathbf{v}) = \lambda \mathbf{a}(\mathbf{u}, \mathbf{v}) \quad \forall \mathbf{u}, \mathbf{v} \in \mathbf{V}, \forall \lambda \in \mathbb{R}$. This holds because you can pull constants out of integrals.
4. $\mathbf{a}(\mathbf{u}, \mathbf{u}) \geq 0 \quad \forall \mathbf{u} \in \mathbf{V}$ This holds because $\mathbf{u}'^2(x) \geq 0$ for all x .
5. $\mathbf{a}(\mathbf{u}, \mathbf{u}) = 0 \iff \mathbf{u} = 0$. It's clear that $\mathbf{u} = 0$ implies $\mathbf{a}(\mathbf{u}, \mathbf{u}) = 0$. We can rigorously prove the other direction as well. Here is a nearly-rigorous proof of why $\mathbf{a}(\mathbf{u}, \mathbf{u}) = 0$ implies $\mathbf{u} = 0$. Fix $x \in [0, 1]$. Then

$$\mathbf{u}(x) = \mathbf{u}(0) + \int_0^x \mathbf{u}'(y) \, dy \quad \implies \quad \mathbf{u}(x) = \int_0^x \mathbf{u}'(y) \, dy.$$

Applying the Schwartz inequality, for $x \in (0, 1]$

$$|u(x)| \leq \left| \int_0^x u'(y) \, dy \right| \leq \sqrt{x} \sqrt{\int_0^x (u'(y))^2 \, dy} \leq \sqrt{a(u, u)}.$$

And so $a(u, u) = 0$ implies $u = 0$ pointwise.