

The Yang-Baxter Technique. Given an algebra U (typically $\hat{U}(\mathfrak{g})$ or $\hat{U}_q(\mathfrak{g})$) and elements $R = \sum a_i \otimes b_i \in U \otimes U$ and $C \in U$, form $Z = \sum_{i,j,k} C a_i b_j a_k C^2 b_i a_j b_k C$.

Problem. Extract information from Z .

The Dogma. Use representation theory. In principle finite, but *slow*.

Definition. A “docile perturbed Gaussian” in the variables $(z_i)_{i \in S}$ over the ring R is an expression of the form

$$\mathbb{e}^{q^{ij} z_i z_j} P = \mathbb{e}^{q^{ij} z_i z_j} \left(\sum_{k \geq 0} \epsilon^k P_k \right),$$

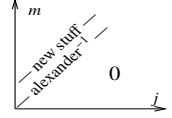
where all coefficients are in R and where P is a “docile series”: $\deg P_k \leq 4k$.

Docility Matters! The rank of the space of docile series to ϵ^k is polynomial in the number of variables $|S|$.

Theorem ([BNG], conjectured [MM], elucidated [Ro1]). Let $J_d(K)$ be the coloured Jones polynomial of K , in the d -dimensional representation of sl_2 . Writing

$$\frac{(q^{1/2} - q^{-1/2}) J_d(K)}{q^{d/2} - q^{-d/2}} \Big|_{q=e^h} = \sum_{j,m \geq 0} a_{jm}(K) d^j h^m,$$

“below diagonal” coefficients vanish, $a_{jm}(K) = 0$ if $j > m$, and “on diagonal” coefficients give the inverse of the Alexander polynomial:



$(\sum_{m=0}^{\infty} a_{mm}(K) h^m) \cdot \omega(K)(e^h) = 1$.

“Above diagonal” we have **Rozansky’s Theorem** [Ro3, (1.2)]:

$$J_d(K)(q) = \frac{q^d - q^{-d}}{(q - q^{-1}) \omega(K)(q^d)} \left(1 + \sum_{k=1}^{\infty} \frac{(q-1)^k \rho_k(K)(q^d)}{\omega^{2k}(K)(q^d)} \right).$$

Prior art. Some amazing computations by Rozansky and Overbay in [Ro2, Ro3] and in [Ov].



Faddeev’s Formula (In as much as we can tell, first appeared w/o proof in Faddeev [Fa], rediscovered and proven in Quesne [Qu], and again with easier proof, in Zagier [Za]). With $[n]_q := \frac{q^n - 1}{q - 1}$, with $[n]_q! := [1]_q [2]_q \cdots [n]_q$ and with $\mathbb{e}_q^x := \sum_{n \geq 0} \frac{x^n}{[n]_q!}$, we have



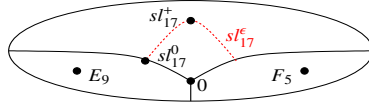
$$\log \mathbb{e}_q^x = \sum_{k \geq 1} \frac{(1-q)^k x^k}{k(1-q^k)} = x + \frac{(1-q)^2 x^2}{2(1-q^2)} + \dots$$

Proof. We have that $\mathbb{e}_q^x = \frac{\mathbb{e}_q^{qx} - \mathbb{e}_q^x}{qx - x}$ (“the q -derivative of \mathbb{e}_q^x is itself”), and hence $\mathbb{e}_q^{qx} = (1 + (1-q)x) \mathbb{e}_q^x$, and

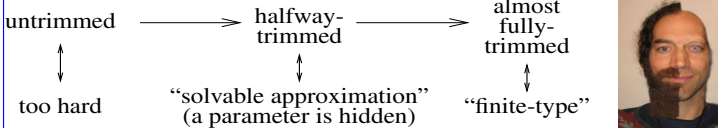
$$\log \mathbb{e}_q^{qx} = \log(1 + (1-q)x) + \log \mathbb{e}_q^x.$$

Writing $\log \mathbb{e}_q^x = \sum_{k \geq 1} a_k x^k$ and comparing powers of x , we get $q^k a_k = -(1-q)^k/k + a_k$, or $a_k = \frac{(1-q)^k}{k(1-q^k)}$. \square

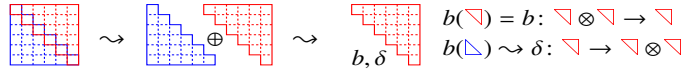
The (fake) moduli of Lie algebras on V , a quadratic variety in $(V^*)^{\otimes 2} \otimes V$ is on the right. We care about $sl_{17}^\epsilon := sl_{17}^\epsilon / (\epsilon^{k+1} = 0)$.



Solvable Approximation. A quantized universal enveloping algebra (aka “quantum group”) is an ∞ -dimensional inverse limit.



Recomposing gl_n . Half is enough! $gl_n \oplus \mathfrak{a}_n = \mathcal{D}(\nabla, b, \delta)$:



Now define $gl_n^\epsilon := \mathcal{D}(\nabla, b, \epsilon\delta)$. Schematically, this is $[\nabla, \nabla] = \nabla$, $[\Delta, \Delta] = \epsilon\Delta$, and $[\nabla, \Delta] = \Delta + \epsilon\nabla$. In detail, it is

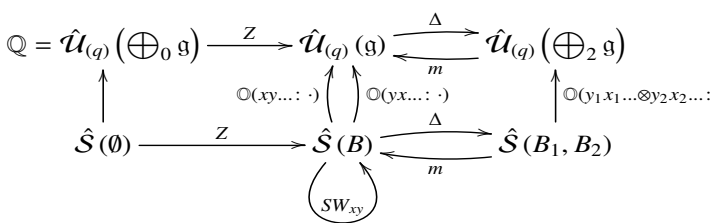
$[e_{ij}, e_{kl}] = \delta_{jk} e_{il} - \delta_{il} e_{kj}$ $[f_{ij}, f_{kl}] = \epsilon \delta_{jk} f_{il} - \epsilon \delta_{il} f_{kj}$
 $[e_{ij}, f_{kl}] = \delta_{jk} (\epsilon \delta_{i < k} e_{il} + \delta_{il} (h_j + \epsilon g_i) / 2 + \delta_{i > l} f_{il})$
 $\quad - \delta_{il} (\epsilon \delta_{k < j} e_{kj} + \delta_{kj} (h_j + \epsilon g_j) / 2 + \delta_{k > j} f_{kj})$
 $[g_i, e_{jk}] = (\delta_{ij} - \delta_{ik}) e_{jk}$ $[h_i, e_{jk}] = \epsilon (\delta_{ij} - \delta_{ik}) e_{jk}$
 $[g_i, f_{jk}] = (\delta_{ij} - \delta_{ik}) f_{jk}$ $[h_i, f_{jk}] = \epsilon (\delta_{ij} - \delta_{ik}) f_{jk}$

Solvable Approximation (2). At $\epsilon = 1$ and modulo $h = g$, the above is just gl_n . By rescaling at $\epsilon \neq 0$, gl_n^ϵ is independent of ϵ . We let gl_n^k be gl_n^ϵ regarded as an algebra over $\mathbb{Q}[\epsilon] / \epsilon^{k+1} = 0$. It is the “ k -smidgen solvable approximation” of gl_n !

Recall that \mathfrak{g} is “solvable” if iterated commutators in it ultimately vanish: $\mathfrak{g}_2 := [\mathfrak{g}, \mathfrak{g}]$, $\mathfrak{g}_3 := [\mathfrak{g}_2, \mathfrak{g}_2]$, \dots , $\mathfrak{g}_d = 0$. Equivalently, if it is a subalgebra of some large-size ∇ algebra.

Note. This whole process makes sense for arbitrary semi-simple Lie algebras.

GDO-Categories. Given \mathfrak{g} with basis $B = \{x, y, \dots\}$, consider the following diagram:



Hence Z , SW_{xy} , m , Δ , (and likewise S and θ) are morphisms in the completion of the monoidal category \mathcal{F} whose objects are finite sets B and whose morphisms are $\text{mor}_{\mathcal{F}}(B, B') := \text{Hom}_{\mathbb{Q}}(\mathcal{S}(B) \rightarrow \mathcal{S}(B')) = \mathcal{S}(B^*, B')$ (by convention, $x^* = \xi$, $y^* = \eta$, etc.). Ergo we need to *consolidate* (at least parts of) said completion.

Aside. “Consolidate” means “give a finite name to an infinite object, and figure out how to sufficiently manipulate such finite names”. E.g., solving $f''' = -f$ we encounter and set $\sum \frac{(-1)^k x^{2k}}{(2k)!} \rightsquigarrow \cos x$, $\sum \frac{(-1)^k x^{2k+1}}{(2k+1)!} \rightsquigarrow \sin x$, and then $\cos^2 x + \sin^2 x = 1$ and $\sin(x+y) = \sin x \cos y + \cos x \sin y$.

The Composition Law. If

$$\mathcal{S}(B_0) \xrightarrow{f} \mathcal{S}(B_1) \xrightarrow{g} \mathcal{S}(B_2)$$

$f \in \mathbb{Q} \llbracket \zeta_{0i}, z_{1j} \rrbracket$ $g \in \mathbb{Q} \llbracket \zeta_{1j}, z_{2k} \rrbracket$

then ${}^t(f \parallel g) = {}^t(g \circ f) = \left(g|_{\zeta_{1j} \rightarrow z_{1j}} f \right)_{z_{1j}=0}$.

Examples.

- The 1-variable identity map $I: \mathcal{S}(z) \rightarrow \mathcal{S}(z)$ is given by ${}^t I_1 = \mathbb{e}^{zz}$ and the n -variable one by ${}^t I_n = \mathbb{e}^{z_1 \zeta_1 + \dots + z_n \zeta_n}$.