

BIG FREQUENCY CASCADES IN THE CUBIC  
NONLINEAR SCHRÖDINGER FLOW ON THE  
2-TORUS

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# 1. INTRODUCTION

# THE NLS INITIAL VALUE PROBLEM

[Joint work with **Keel, Staffilani, Takaoka and Tao**]

We consider the defocusing initial value problem:

$$\begin{cases} (-i\partial_t + \Delta)u = |u|^2 u \\ u(0, x) = u_0(x), \text{ where } x \in \mathbb{T}^2. \end{cases} \quad (\text{NLS}(\mathbb{T}^2))$$

Smooth solution  $u(x, t)$  exists globally and

$$\text{Mass} = M(u) = \|u(t)\|^2 = M(0)$$

$$\text{Energy} = E(u) = \int \left( \frac{1}{2} |\nabla u(t, x)|^2 + \frac{1}{4} |u(x, t)|^4 \right) dx = E(0)$$

We want to understand the shape of  $|\hat{u}(t, \xi)|$ . The conservation laws impose  $L^2$ -moment constraints on this object.

# NOTION OF FREQUENCY CASCADE

## DEFINITION

A frequency cascade is the phenomenon of global-in-time solutions shifting their mass toward increasingly high frequencies.

This shift is also called a **forward cascade**.

- A way to measure the cascade is to study

$$\|u(t)\|_{\dot{H}^s}^2 = \int |\hat{u}(t, \xi)|^2 |\xi|^{2s} d\xi$$

and prove that it grows for large times  $t$ .

- The cascade is incompatible with **scattering** and **integrability**.

# INCOMPATIBLE WITH SCATTERING & INTEGRABILITY

- **Scattering:**  $\forall$  global solution  $u(t, x) \in H^s \exists u_0^+ \in H^s$  such that,

$$\lim_{t \rightarrow +\infty} \|u(t, x) - e^{it\Delta} u_0^+(x)\|_{H^s} = 0.$$

Note:  $\|e^{it\Delta} u_0^+\|_{H^s} = \|u_0^+\|_{H^s} \implies \|u(t)\|_{H^s}$  is bounded.

- **Complete Integrability:** The 1d equation

$$(i\partial_t + \Delta)u = -|u|^2 u$$

has infinitely many conservation laws. Combining them in the right way one gets that  $\|u(t)\|_{H^s} \leq C_s$  for all times.

## PAST RESULTS

- Bourgain: (late 90's)

For the periodic IVP  $NLS(\mathbb{T}^2)$  one can prove

$$\|u(t)\|_{H^s}^2 \leq C_s |t|^{4s}.$$

The idea is to improve the local estimate for  $t \in [-1, 1]$

$$\|u(t)\|_{H^s} \leq C_s \|u(0)\|_{H^s}, \quad \text{for } C_s \gg 1$$

( $\implies \|u(t)\|_{H^s} \lesssim C^{|t|}$  upper bounds) to obtain

$$\|u(t)\|_{H^s} \leq 1 \|u(0)\|_{H^s} + C_s \|u(0)\|_{H^s}^{1-\delta} \quad \text{for } C_s \gg 1,$$

for some  $\delta > 0$ . This iterates to give

$$\|u(t)\|_{H^s} \leq C_s |t|^{1/\delta}.$$

- Improvements: Staffilani, Colliander-Delort-Kenig-Staffilani.

## PAST RESULTS

- Bourgain: (late 90's)

Given  $m, s \gg 1$  there exist  $\tilde{\Delta}$  and a global solution  $u(x, t)$  to the modified wave equation

$$(\partial_{tt} - \tilde{\Delta})u = u^p$$

such that  $\|u(t)\|_{H^s} \sim |t|^m$ .

- Physics: Weak turbulence theory: Hasselmann & Zakharov.  
Numerics (d=1): Majda-McLaughlin-Tabak; Zakharov et. al.
- Kuksin has studied a small dispersion NLS

$$i\partial_t w + \delta\Delta w = |w|^2 w$$

with odd, periodic boundary conditions and with  $0 < \delta \ll 1$ . Smooth norms of relatively generic data, with unit sized  $L^2$  norm, are shown to grow larger than a negative power of  $\delta$ . These results correspond to large data solutions of the  $\delta = 1$  problem  $NLS_3^+(\mathbb{T}^2)$ .



# EXPLODING SOBOLEV NORMS CONJECTURE?

- Solutions to dispersive equations on  $\mathbb{R}^d$  have bounded high Sobolev norms.
- There are solutions to nonlinear dispersive equations on  $\mathbb{T}^d$  with exploding Sobolev norms. In particular for  $NLS(\mathbb{T}^2)$  there exists  $u(t, x)$  such that

$$\|u(t)\|_{H^s}^2 \rightarrow \infty \text{ as } t \rightarrow \infty.$$

- The cascade phenomena of high Sobolev norm explosion should be generic in phase space.
- Z. Hani developed a strategy (and and proved some partial results) toward proving existence and genericity of infinite cascades.

# MAIN RESULT

We consider the defocusing initial value problem:

$$\begin{cases} (-i\partial_t + \Delta)u = |u|^2 u \\ u(0, x) = u_0(x), \text{ where } x \in \mathbb{T}^2. \end{cases} \quad (\text{NLS}(\mathbb{T}^2))$$

## THEOREM (C-KEEL-STAFFILANI-TAKAOKA-TAO)

Let  $s > 1$ ,  $K \gg 1$  and  $0 < \sigma < 1$  be given. Then there exists a global smooth solution  $u(t, x)$  and  $T > 0$  such that

$$\|u_0\|_{H^s} \leq \sigma$$

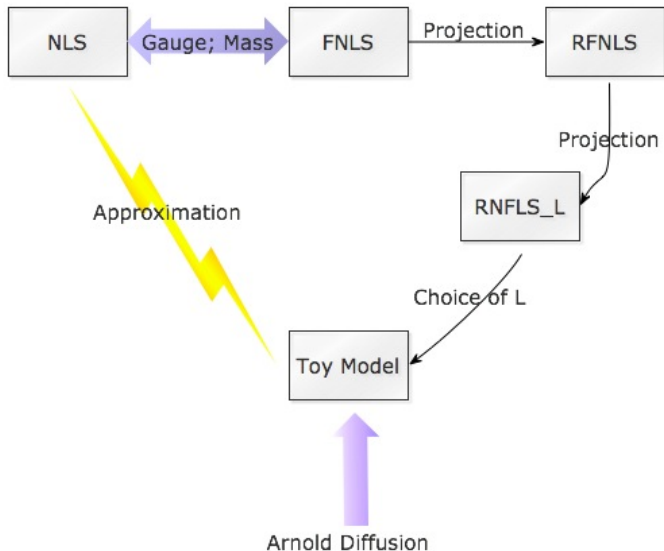
and

$$\|u(t)\|_{H^s}^2 \geq K.$$

M. Guardia and V. Kaloshin (preprint 2012 ): polynomial speed!

## 2. OVERVIEW OF PROOF

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# PRELIMINARY REDUCTIONS

- **Gauge Freedom:**

If  $u$  solves NLS then  $v(t, x) = e^{-i2Gt} u(t, x)$  solves

$$\begin{cases} i\partial_t v + \Delta v = (2G + |v|^2)v \\ v(0, x) = v_0(x), \end{cases} \quad x \in \mathbb{T}^2. \quad (\text{NLS}_G)$$

- **Fourier Ansatz:** Recast the dynamics in Fourier coefficients,

$$v(t, x) = \sum_{n \in \mathbb{Z}^2} a_n(t) e^{i(n \cdot x + |n|^2 t)}.$$

$$\begin{cases} i\partial_t a_n = 2Ga_n + \sum_{\substack{n_1, n_2, n_3 \in \mathbb{Z}^2 \\ n_1 - n_2 + n_3 = n}} a_{n_1} \bar{a}_{n_2} a_{n_3} e^{i\omega_4 t} \\ a_n(0) = \hat{u}_0(n), \end{cases} \quad n \in \mathbb{Z}^2. \quad (\mathcal{FNLS}_G)$$

# PRELIMINARY REDUCTIONS

## ■ Diagonal decomposition of sum:

$$\begin{aligned} \sum_{\substack{n_1, n_2, n_3 \in \mathbb{Z}^2 \\ n_1 - n_2 + n_3 = n}} &= \sum_{\substack{n_1, n_2, n_3 \in \mathbb{Z}^2 \\ n_1 - n_2 + n_3 = n \\ n \neq n_1, n_3}} + \sum_{\substack{n_1, n_2, n_3 \in \mathbb{Z}^2 \\ n_1 - n_2 + n_3 = n \\ n = n_1}} \\ &+ \sum_{\substack{n_1, n_2, n_3 \in \mathbb{Z}^2 \\ n_1 - n_2 + n_3 = n \\ n = n_3}} - \sum_{\substack{n_1, n_2, n_3 \in \mathbb{Z}^2 \\ n_1 - n_2 + n_3 = n \\ n = n_1 = n_3}} \end{aligned}$$

## ■ Choice of $G$ :

$$G = -\|u_0\|_{L^2}^2.$$

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# RESONANT TRUNCATION

- NLS dynamic is recast as

$$-i\partial_t a_n = -a_n |a_n|^2 + \sum_{n_1, n_2, n_3 \in \Gamma(n)} a_{n_1} \bar{a}_{n_2} a_{n_3} e^{i\omega_4 t}. \quad (\mathcal{FNLS})$$

where  $\omega_4 = |n_1|^2 - |n_2|^2 + |n_3|^2 - |n|^2$ , and

$$\Gamma(n) = \{n_1, n_2, n_3 \in \mathbb{Z}^2 : n_1 - n_2 + n_3 = n, n_1 \neq n, n_3 \neq n\}.$$

- 

$$\begin{aligned} \Gamma_{res}(n) &= \{n_1, n_2, n_3 \in \Gamma(n) : \omega_4 = 0\}. \\ &= \{ \text{Triples } (n_1, n_2, n_3) : (n_1, n_2, n_3, n) \text{ is a rectangle} \} \end{aligned}$$

- The resonant truncation of  $\mathcal{FNLS}$  is

$$-i\partial_t b_n = -b_n |b_n|^2 + \sum_{n_1, n_2, n_3 \in \Gamma_{res}(n)} b_{n_1} \bar{b}_{n_2} b_{n_3}. \quad (R\mathcal{FNLS})$$



# FINITE DIMENSIONAL RESONANT RESTRICTION

- A set  $\Lambda \subset \mathbb{Z}^2$  is *closed under resonant interactions* if

$$n_1, n_2, n_3 \in \Gamma_{res}(n), n_1, n_2, n_3 \in \Lambda \implies n \in \Lambda.$$

- A *finite dimensional resonant restriction* of  $\mathcal{FNLS}$  is

$$-i\partial_t b_n = -b_n |b_n|^2 + \sum_{n_1, n_2, n_3 \in \Gamma_{res}(n) \cap \Lambda^3} b_{n_1} \bar{b}_{n_2} b_{n_3}. \quad (R\mathcal{FNLS}_\Lambda)$$

- $\forall$  resonant-closed finite  $\Lambda \subset \mathbb{Z}^2$   $R\mathcal{FNLS}_\Lambda$  is an ODE.
- If  $\text{spt}(a_n(0)) \subset \Lambda$  then  $\mathcal{FNLS}$ -evolution  $a_n(0) \mapsto a_n(t)$  is nicely approximated by  $R\mathcal{FNLS}_\Lambda$ -ODE  $a_n(0) \mapsto b_n(t)$ .
- Given  $\epsilon, s, K$ , build  $\Lambda$  so that  $R\mathcal{FNLS}_\Lambda$  cascades.

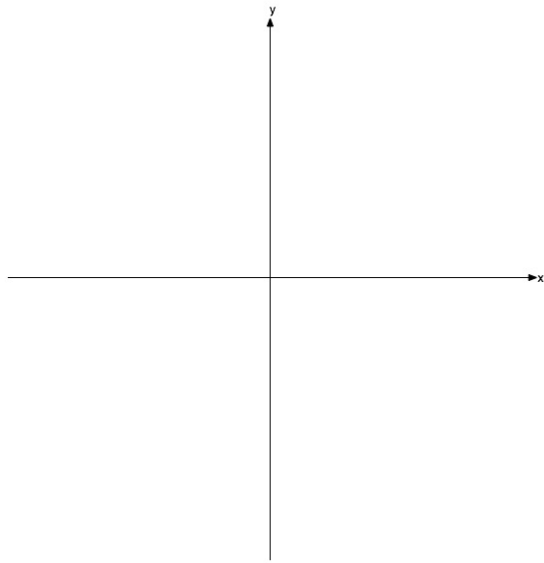
# IMAGINE WE BUILD A RESONANT $\Lambda \subset \mathbb{Z}^2$ SUCH THAT...

Imagine a resonant-closed  $\Lambda = \Lambda_1 \cup \dots \cup \Lambda_M$  with **properties**.

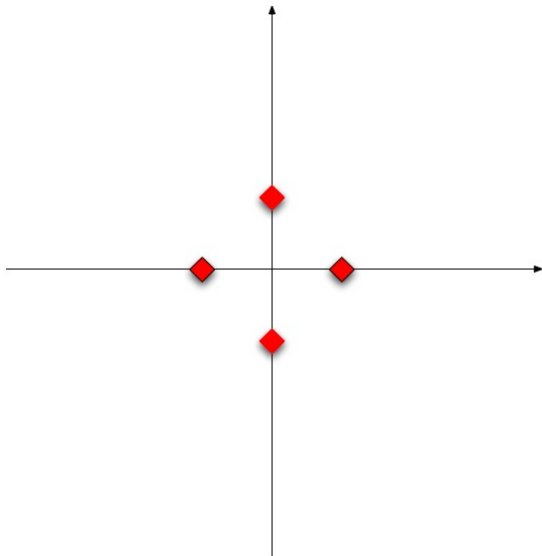
Define a **nuclear family** to be a rectangle  $(n_1, n_2, n_3, n_4)$  where the frequencies  $n_1, n_3$  (the 'parents') live in **generation**  $\Lambda_j$  and  $n_2, n_4$  ('children') live in generation  $\Lambda_{j+1}$ .

- $\forall 1 \leq j < M$  and  $\forall n_1 \in \Lambda_j \exists$  unique nuclear family such that  $n_1, n_3 \in \Lambda_j$  are parents and  $n_2, n_4 \in \Lambda_{j+1}$  are children.
- $\forall 1 \leq j < M$  and  $\forall n_2 \in \Lambda_{j+1} \exists$  unique nuclear family such that  $n_2, n_4 \in \Lambda_{j+1}$  are children and  $n_1, n_3 \in \Lambda_j$  are parents.
- The sibling of a frequency is never its spouse.
- Besides nuclear families,  $\Lambda$  contains no other rectangles.
- The function  $n \mapsto a_n(0)$  is constant on each generation  $\Lambda_j$ .

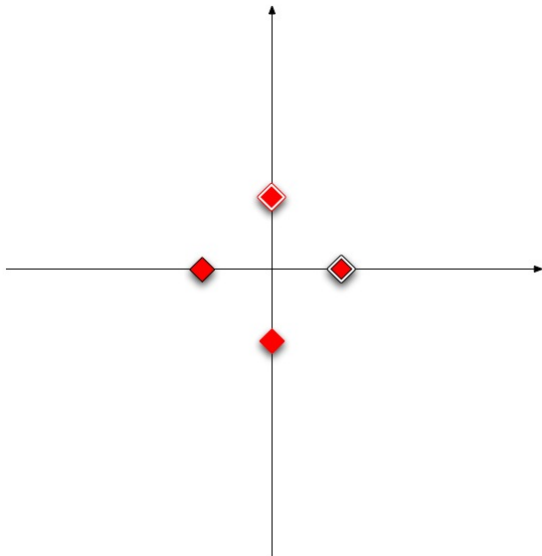
# CARTOON CONSTRUCTION OF $\Lambda$



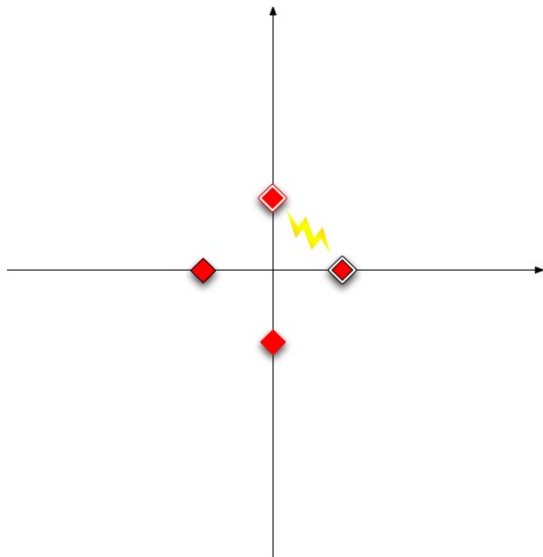
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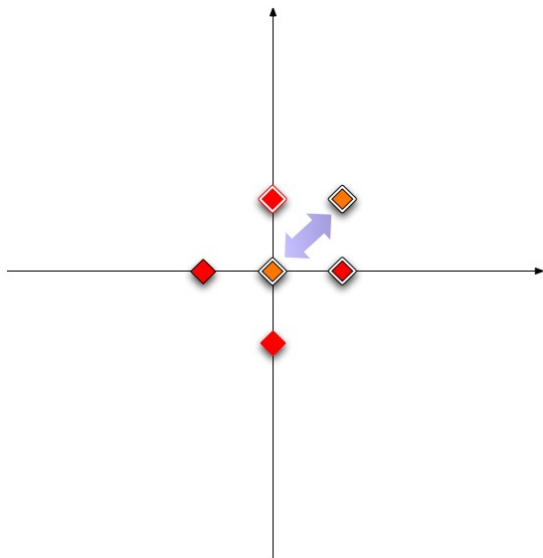
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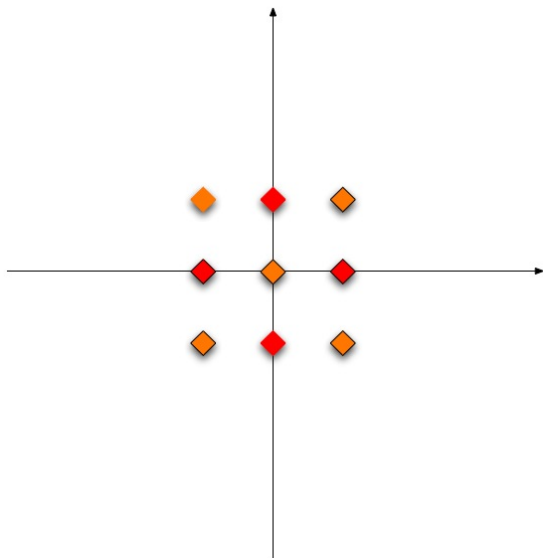
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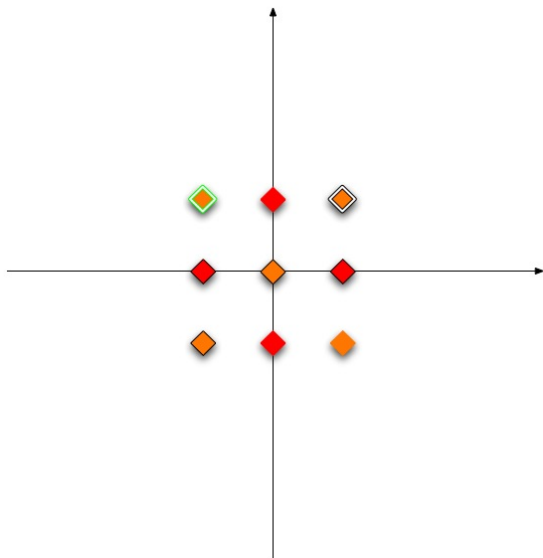


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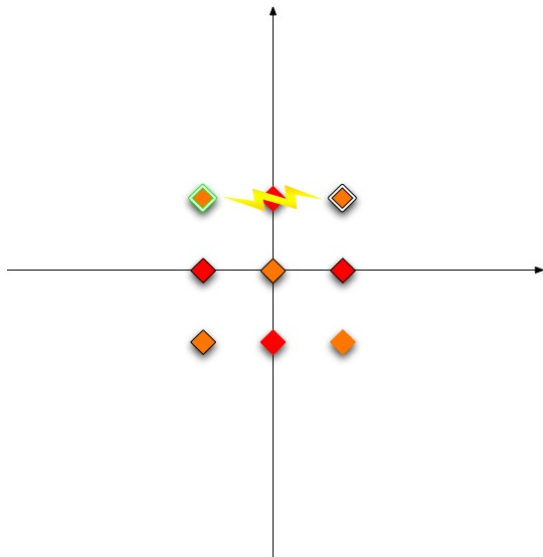




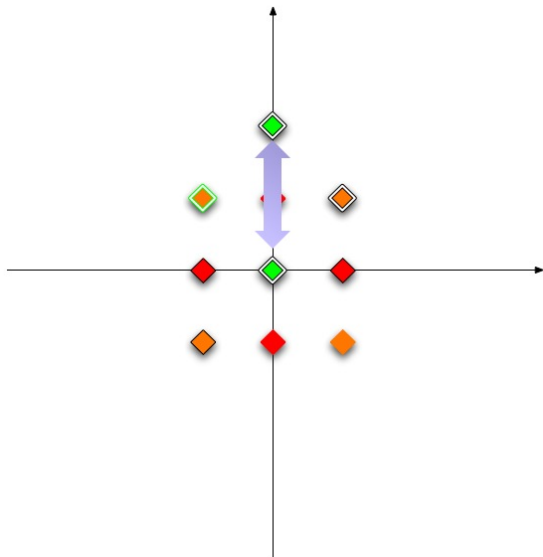
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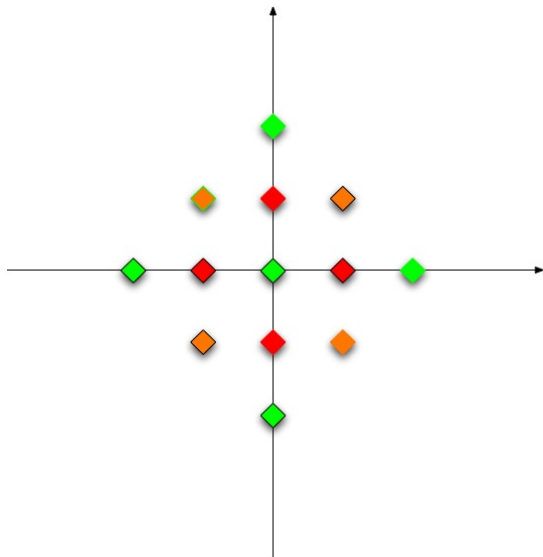
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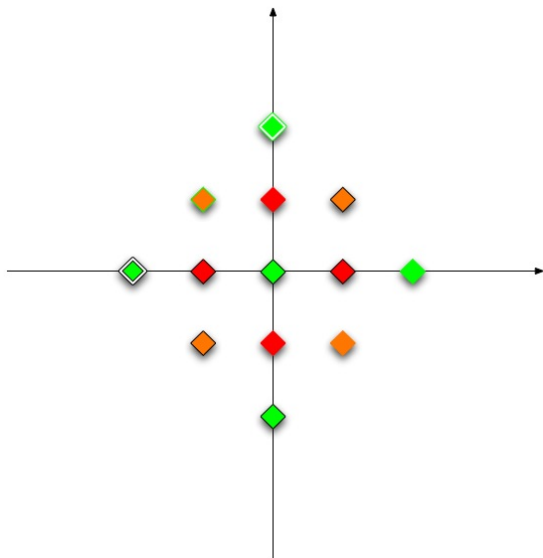
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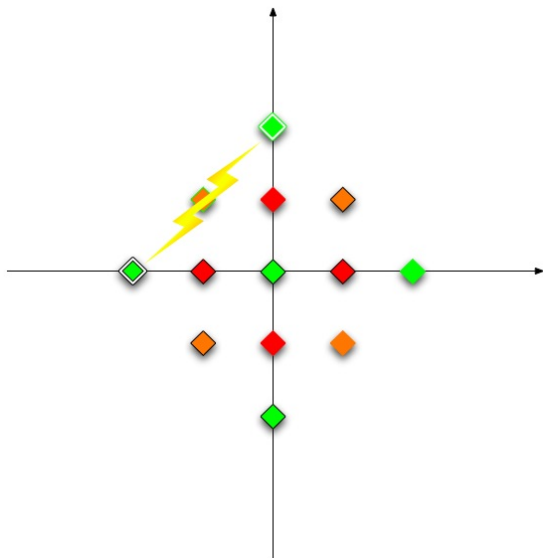
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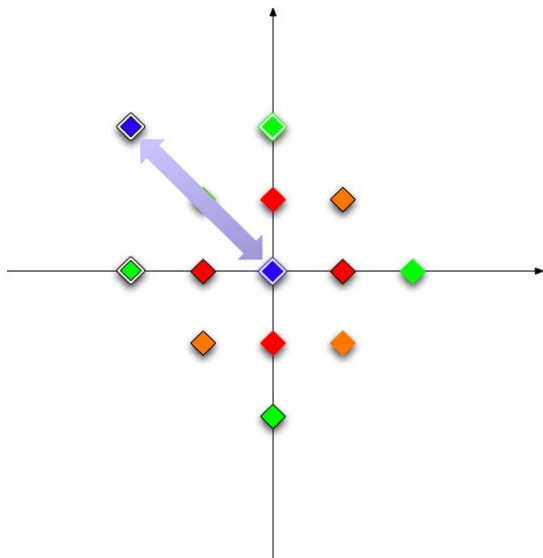
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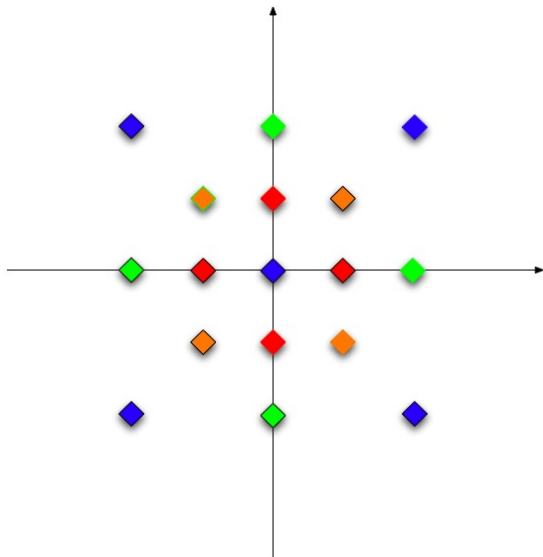
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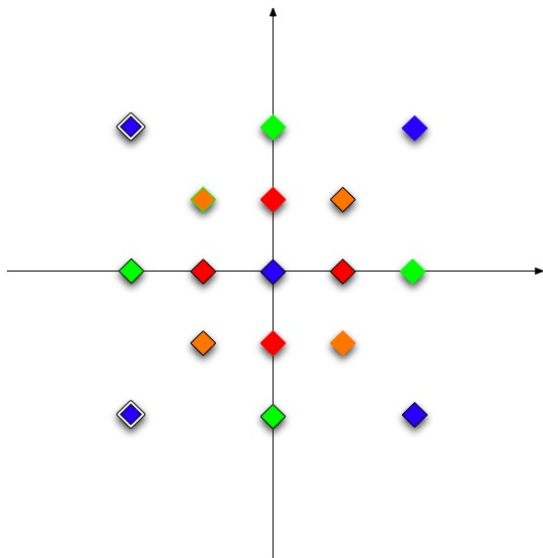


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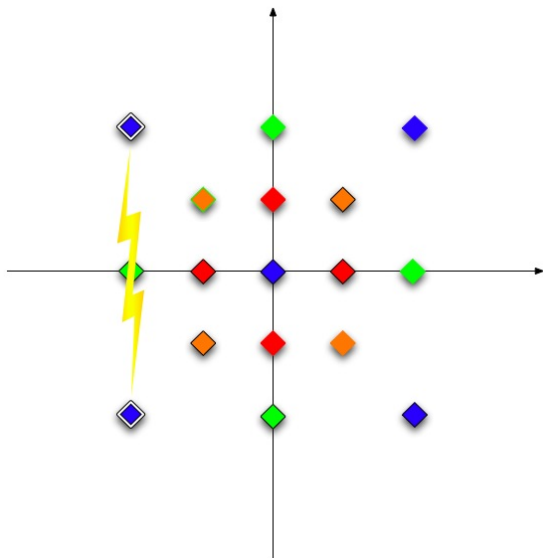




# CARTOON CONSTRUCTION OF $\Lambda$



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# THE TOY MODEL ODE

**Assume** we can construct such a  $\Lambda = \Lambda_1 \cup \dots \cup \Lambda_M$ . The properties imply  $R\mathcal{FNLS}_\Lambda$  simplifies to the **toy model** ODE

$$\partial_t b_j(t) = -i|b_j(t)|^2 b_j(t) + 2i\bar{b}_j(t)[b_{j-1}(t)^2 + b_{j+1}(t)^2].$$

$$L^2 \sim \sum_j |b_j(t)|^2 = \sum_j |b_j(0)|^2$$

$$H^s \sim \sum_j |b_j(t)|^2 \left( \sum_{n \in \Lambda_j} |n|^{2s} \right).$$

We also want  $\Lambda$  to satisfy **Wide Diaspora Property**

$$\sum_{n \in \Lambda_M} |n|^{2s} \gg \sum_{n \in \Lambda_1} |n|^{2s}.$$

# CONSERVATION LAWS FOR THE *ODE* SYSTEM

$$\text{Mass} = \sum_j |b_j(t)|^2 = C_0$$

$$\text{Momentum} = \sum_j |b_j(t)|^2 \sum_{n \in \Lambda_j} n = C_1,$$

$$\text{Energy} = K + P = C_2,$$

where

$$K = \sum_j |b_j(t)|^2 \sum_{n \in \Lambda_j} |n|^2,$$

$$P = \frac{1}{2} \sum_j |b_j(t)|^4 + \sum_j |b_j(t)|^2 |b_{j+1}(t)|^2.$$

**Conservation laws for ODE do not involve Fourier moments!**

### 3. CONCATENATED SLIDERS FOR TOY MODEL ODE

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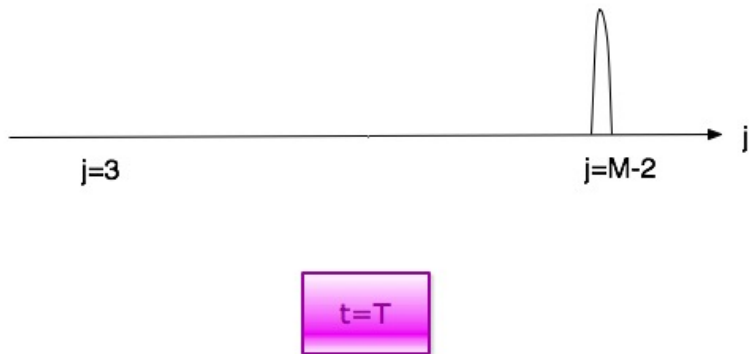
Using dynamical systems methods, we construct a Toy Model ODE evolution such that:



$t=0$

### 3. CONCATENATED SLIDERS FOR TOY MODEL ODE

Using dynamical systems methods, we construct a Toy Model ODE evolution such that:



A travelling wave through the generations.

# PROPERTIES OF THE TOY MODEL *ODE*

- Solution of the Toy Model is a vector flow  $t \rightarrow b(t) \in \mathbb{C}^M$   
$$b(t) = (b_1(t), \dots, b_M(t)) \in \mathbb{C}^M; b_j = 0 \forall j \leq 0, j \geq M + 1.$$
- Local Well-Posedness; Let  $S(t)$  denote associated flowmap.
- Mass Conservation:  $|b(t)|^2 = |b(0)|^2 \implies$ 
  - Toy Model ODE is Globally Well-Posed.
  - Invariance of the sphere:  $\Sigma = \{x \in \mathbb{C}^M : |x|^2 = 1\}$

$$S(t)\Sigma = \Sigma.$$



# PROPERTIES OF THE TOY MODEL *ODE*

- Support Conservation:

$$\begin{aligned}\partial_t |b_j|^2 &= 2\operatorname{Re}(\overline{b_j} \partial_t b_j) \\ &= 4\operatorname{Re}(i\overline{b_j}^2 [b_{j-1}^2 + b_{j+1}^2]) \\ &\leq 4|b_j|^2.\end{aligned}$$

Thus, if  $b_j(0) = 0$  then  $b_j(t) = 0$  for all  $t$ .

- Invariance of coordinate tori:

$$\mathbb{T}_j = \{(b_1, \dots, b_M \in \Sigma) : |b_j| = 1, b_k = 0 \forall k \neq j\}$$

Mass Conservation  $\implies S(T)\mathbb{T}_j = \mathbb{T}_j$ .

Dynamics on the invariant tori is easy:

$$b_j(t) = e^{-i(t+\theta)}; b_k(t) = 0 \forall k \neq j.$$

# EXPLICIT SLIDER SOLUTIONS

Consider  $M = 2$ . Then *ODE* is of the form

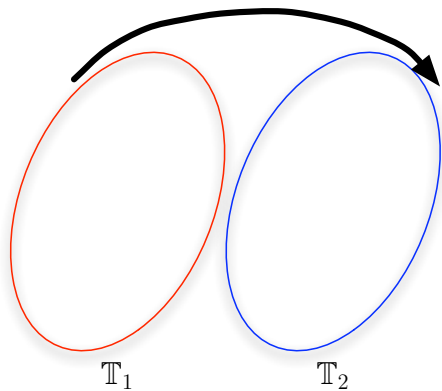
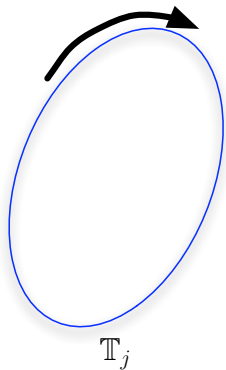
$$\begin{aligned}\partial_t b_1 &= -i|b_1|^2 b_1 + 2i\bar{b}_1 b_2^2 \\ \partial_t b_2 &= -i|b_2|^2 b_2 + 2i\bar{b}_2 b_1^2.\end{aligned}$$

Let  $\omega = e^{2i\pi/3}$  (cube root of unity). This ODE has explicit solution

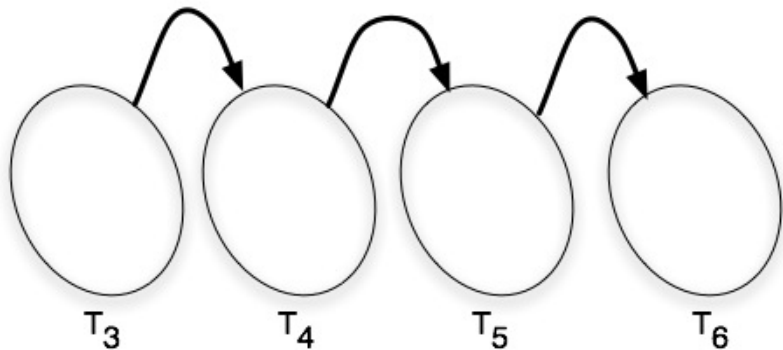
$$b_1(t) = \frac{e^{-it}}{\sqrt{1 + e^{2\sqrt{3}t}}}\omega, \quad b_2(t) = \frac{e^{-it}}{\sqrt{1 + e^{-2\sqrt{3}t}}}\omega^2.$$

- As  $t \rightarrow -\infty$ ,  $(b_1(t), b_2(t)) \rightarrow (e^{-it}\omega, 0) \in \mathbb{T}_1$ .
- As  $t \rightarrow +\infty$ ,  $(b_1(t), b_2(t)) \rightarrow (0, e^{-it}\omega^2) \in \mathbb{T}_2$ .

# TWO EXPLICIT SOLUTION FAMILIES



## CONCATENATED SLIDERS: IDEA OF PROOF



# THE PERFECT SHOT?

*Off the expressway, over the river, off the billboard, through the window, nothin but net. —Michael Jordan*

# DIFFUSION FOR TOY MODEL STATEMENT

## THEOREM

*Let  $M \geq 6$ . Given  $\epsilon > 0$  there exist  $x_3$  within  $\epsilon$  of  $\mathbb{T}_3$  and  $x_{M-2}$  within  $\epsilon$  of  $\mathbb{T}_{M-2}$  and a time  $t$  such that*

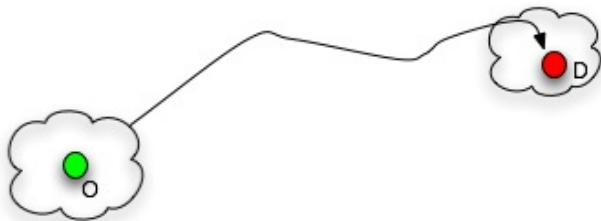
$$S(t)x_3 = x_{M-2}.$$

## REMARK

*$S(t)x_3$  is a solution of total mass 1 arbitrarily concentrated at mode  $j = 3$  at some time  $t_0$  and then arbitrarily concentrated at mode  $j = M - 2$  at later time  $t$ .*

# TARGETS AND COVERING

Let  $O, D$  denote points in our phase space  $\Sigma$ . Can we flow along  $S(t)$  from *nearby* the origin point  $O$  to *nearby* the destination point  $D$ ? More generally, suppose  $O$  and  $D$  are subsets of  $\Sigma$ .



The notion of a **target** quantifies this question.

# TARGETS

- Let  $\mathcal{M}$  denote a subset of  $\Sigma$ . Let  $d$  be a (pseudo)metric on  $\Sigma$ . Let  $R > 0$  be a radius.
- The **Target**  $(\mathcal{M}, d, R) := \{x \in \Sigma : d(x, \mathcal{M}) < R\}$ .
- Given  $x, y \in \Sigma$ .  
We say  $x$  **hits**  $y$  if  $y = S(t)x$  for some  $t \geq 0$ .



# COVERING

Given an initial target  $(M_1, d_1, R_1)$  and a final target  $(M_2, d_2, R_2)$ . We say  $(M_1, d_1, R_1)$  **can cover**  $(M_2, d_2, R_2)$  and write

$$(M_1, d_1, R_1) \implies (M_2, d_2, R_2)$$

if:

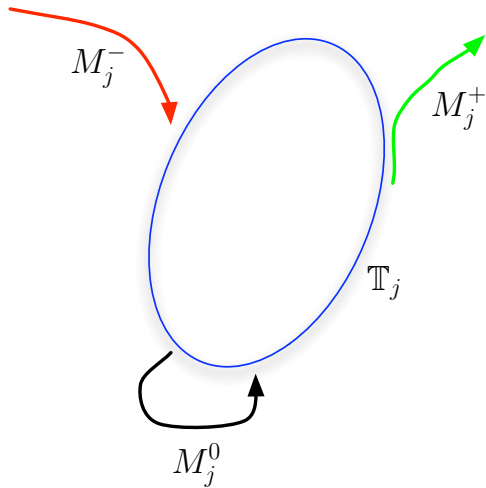
$\forall x_2 \in M_2 \exists x_1 \in M_1$  such that  $\forall y_1 \in \Sigma$  with  $d_1(x_1, y_1) < R_1 \exists y_2 \in \Sigma$  with  $d(x_2, y_2) < R_2$  such that  $y_1$  hits  $y_2$ .

- The flowout of  $(M_1, d_1, R_1)$  is **surjective** onto  $(M_2, d_2, R_2)$ .
- Covering also includes a notion of stability.

# STRATEGY OF PROOF

- **Transitivity of Covering:** If  $(M_1, d_1, r_1) \implies (M_2, d_2, r_2)$   
and  $(M_2, d_2, r_2) \implies (M_3, d_3, r_3)$   
then  $(M_1, d_1, r_1) \implies (M_3, d_3, r_3)$ .
- $\forall j \in 3, \dots, M - 2$  we define 3 targets close to  $\mathbb{T}_j$ :
  - Incoming Target  $(M_j^-, d_j^-, R_j^-)$
  - Ricochet Target  $(M_j^0, d_j^0, R_j^0)$
  - Outgoing Target  $(M_j^+, d_j^+, R_j^+)$
- $\forall j = 3, \dots, M - 2$  with appropriate  $d_j^{-,0,+}, R_j^{-,0,+}$ , prove:
  - $(M_j^-, d_j^-, R_j^-) \implies (M_j^0, d_j^0, R_j^0)$
  - $(M_j^0, d_j^0, R_j^0) \implies (M_j^+, d_j^+, R_j^+)$
  - $(M_j^+, d_j^+, R_j^+) \implies (M_{j+1}^-, d_{j+1}^-, R_{j+1}^-)$

# TARGETS AROUND $\mathbb{T}_j$



## 4. CONSTRUCTION OF RESONANT SET $\Lambda$

## 4. CONSTRUCTION OF RESONANT SET $\Lambda$

The task is to **construct a finite set**  $\Lambda \subset \mathbb{Z}^2$  satisfying the **properties** that led to the Toy Model ODE. We do this in two steps:

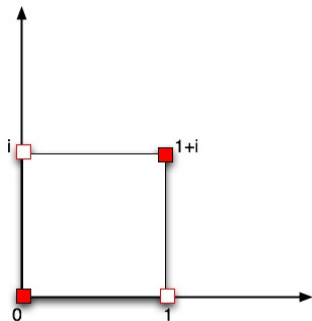
- 1 Build combinatorial model of  $\Lambda$  called  $\Sigma \subset \mathbb{C}^{M-1}$ .
- 2 Build a map  $f : \mathbb{C}^{M-1} \rightarrow \mathbb{R}^2$  which gives

$$f(\Sigma) = \Lambda \subset \mathbb{Z}^2$$

satisfying the **properties**.

# CONSTRUCTION OF COMBINATORIAL MODEL $\Sigma$

- Standard Unit Square:  $S = \{0, 1, 1 + i, i\} \subset \mathbb{C}$ ,  $S = S_1 \cup S_2$  where  $S_1 = \{1, i\}$  and  $S_2 = \{0, 1 + i\}$



- $\mathbb{Z}^2 \equiv \mathbb{Z}[i]; (n_1, n_2) \equiv n_1 + in_2$

# CONSTRUCTION OF COMBINATORIAL MODEL $\Sigma$

- We define

$$\Sigma_j = \{(z_1, z_2, \dots, z_{M-1}) : z_1, \dots, z_{j-1} \in S_2, z_j, \dots, z_{M-1} \in S_1\}$$

with the properties

- $\Sigma_j = S_2^{j-1} \times S_1^{M-j} \subset \mathbb{C}^{M-1}$
  - $|\Sigma_j| = 2^{M-1}$
- Next, we define

$$\Sigma = \Sigma_1 \cup \dots \cup \Sigma_M.$$

- $|\Sigma| = M2^{M-1}$ .
- $\Sigma_j$  is called a generation.

# COMBINATORIAL NUCLEAR FAMILY

- Consider the set  $F = \{F_0, F_1, F_{1+i}, F_i\} \subset \Sigma$  defined by

$$F_w = (z_1, \dots, z_{j-1}, w, z_{j+1}, \dots, z_n)$$

with  $z_1, \dots, z_{j-1} \in S_2$  and  $z_{j+1}, \dots, z_n \in S_2$  and  $w \in S$ .

- The elements  $F_0, F_{1+i} \in \Sigma_{j+1}$  are called *children*.
- The elements  $F_1, F_i$  are called *parents*.
- The four element set  $F$  is called a **combinatorial nuclear family connecting the generations  $\Sigma_j$  and  $\Sigma_{j+1}$** .
- $\forall j \exists 2^{M-2}$  combinatorial nuclear families connecting generations  $\Sigma_j$  and  $\Sigma_{j+1}$ .
- The set  $\Sigma$  satisfies
  - Existence and uniqueness of spouse and children (of sibling and parents).
  - Sibling is never also a spouse.



# CONSTRUCTION OF THE PLACEMENT FUNCTION

We need to map  $\Sigma \subset \mathbb{C}^{M-1}$  into the frequency lattice  $\mathbb{Z}^2$ .

- We first define  $f_1 : \Sigma_1 \rightarrow \mathbb{C}$ .
- $\forall 1 \leq j \leq M$  and each combinatorial nuclear family  $F$  connecting generations  $\Sigma_j$  and  $\Sigma_{j+1}$ , we associate an angle  $\theta(F) \in \mathbb{R}/2\pi\mathbb{Z}$ .
- Given  $f_1$  and the angles of all the families, we define placement functions  $f_j : \Sigma_j \rightarrow \mathbb{C}$  recursively by the rule: Suppose  $f_j : \Sigma_j \rightarrow \mathbb{C}$  has been defined. We define  $f_{j+1} : \Sigma_{j+1} \rightarrow \mathbb{C}$ :

$$f_{j+1}(F_{1+i}) = \frac{1 + e^{i\theta(F)}}{2} f_j(F_1) + \frac{1 - e^{i\theta(F)}}{2} f_j(F_i)$$
$$f_{j+1}(F_0) = \frac{1 + e^{i\theta(F)}}{2} f_j(F_1) - \frac{1 - e^{i\theta(F)}}{2} f_j(F_i)$$

for all combinatorial nuclear families connecting  $\Sigma_j$  to  $\Sigma_{j+1}$ .

# THEOREM: GOOD PLACEMENT FUNCTION

Let  $M \geq 2$ ,  $s > 1$ , and let  $N$  be a sufficiently large integer (depending on  $M$ ).  $\exists$  an initial placement function  $f_1 : \Sigma_1 \rightarrow \mathbb{C}$  and choices of angles  $\theta(F)$  for each nuclear family  $F$  (and thus an associated complete placement function  $f : \Sigma \rightarrow \mathbb{C}$ ) with the following properties:

- **(Non-degeneracy)** The function  $f$  is injective.
- **(Integrality)** We have  $f(\Sigma) \subset \mathbb{Z}[i]$ .
- **(Magnitude)** We have  $C(M)^{-1}N \leq |f(x)| \leq C(M)N$  for all  $x \in \Sigma$ .
- **(Closure/Faithfulness)** If  $x_1, x_2, x_3$  are distinct elements of  $\Sigma$  are such that  $f(x_1), f(x_2), f(x_3)$  form a right-angled triangle, then  $x_1, x_2, x_3$  belong to a combinatorial nuclear family.
- **(Wide Diaspora/Norm Explosion)** We have

$$\sum_{n \in f(\Sigma_M)} |n|^{2s} > \frac{1}{2} 2^{(s-1)(M-1)} \sum_{n \in f(\Sigma_1)} |n|^{2s}.$$