

RECENT PROGRESS ON BLOWUP PHENOMENA IN NONLINEAR SCHRÖDINGER EQUATIONS

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Time Invariant Quantities

$$\text{Mass} = \|u(t)\|_{L_x^2}$$

$$\text{Hamiltonian} = \int_{\mathbb{R}^d} |\nabla u(t)|^2 dx \mp \frac{2}{p+1} |u(t)|^{p+1} dx$$

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- If u solves $NLS_p^\pm(\mathbb{R}^d)$ on $(-T_*, T^*) \times \mathbb{R}^2$ then

$$u_\lambda(\tau, y) := \lambda^{\frac{2}{1-p}} u(\tau \lambda^{-2}, y \lambda^{-1})$$

solves $NLS_p^\pm(\mathbb{R}^d)$ on $(-\lambda^2 T_*, \lambda^2 T^*) \times \mathbb{R}^2$.

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- L^2 and \dot{H}^1 critical cases distinguished by conservation laws.

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- Optimal local-in-time well-posedness (**LWP**) for $NLS_p^\pm(\mathbb{R}^d)$:
 $\forall s \geq \max(0, s_c) \exists$ unique continuous data-to-solution map

$$H^s \ni u_0 \longmapsto u \in C([0, T_{lwp}]; H^s) \cap L_t^q L_x^p$$

with $T_{lwp} = T_{lwp}(\|u_0\|_{H^s})$ if $s > s_c$ and $T_{lwp} = T(u_0)$ if $s = s_c$.

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- Optimal maximal-in-time well-posedness (**GWP**) is known only in the defocusing energy critical case. What is the fate of local-in-time solutions with critical initial regularity?

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- $\forall u_0 \in L^2$ there exists $T_{lwp}(u_0)$ determined by

$$\|e^{it\Delta} u_0\|_{L_{tx}^4([0, T_{lwp}] \times \mathbb{R}^2)} < \frac{1}{100}.$$

\exists unique solution $u \in C([0, T_{lwp}]; L^2) \cap L_{tx}^4([0, T_{lwp}] \times \mathbb{R}^2)$.

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- Define the **maximal forward existence time** $T^*(u_0)$ by

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for all $\delta > 0$ but diverges to ∞ as $\delta \downarrow 0$.

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- \exists **small data scattering threshold** $\mu_0 > 0$

$$\|u_0\|_{L^2} < \mu_0 \implies \|u\|_{L_{tx}^4(\mathbb{R} \times \mathbb{R}^2)} < 2\mu_0.$$

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$NLS_5^+(\mathbb{R}^1)$ is similarly H^s -GWP for $s > \frac{4}{9}$.

[Tzirakis]

$NLS_{\frac{4}{d}+1}^+(\mathbb{R}^d)$ is H^s -GWP for $s > \frac{d+8}{d+10}$.

[Visan-Zhang]

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- Arise as *pseudoconformal* image of $e^{it}Q(x)$:

$$S(t, x) = \frac{1}{t} Q\left(\frac{x}{t}\right) e^{-i\frac{|x|^2}{4t} + \frac{i}{t}}.$$

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$$\|S(-1)\|_{L_x^2} = \|Q\|_{L^2}.$$

All mass in S is *conically* concentrated into a point.

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- **Minimal mass H^1 blowup solution characterization:**
 $u_0 \in H^1$, $\|u_0\|_{L^2} = \|Q\|_{L^2}$, $T^*(u_0) < \infty$ implies that $u = S$ up to an explicit solution symmetry. [Merle]

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- How do these solutions blow up?

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$$\liminf_{t \uparrow T^*} \int_{|x| < (T^* - t)^{1/2 -}} |u(t, x)|^2 dx \geq \|Q\|_{L^2}^2.$$

[Merle-Tsutsumi]

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- Fantastic recent progress on the H^1 blowup theory.

[Merle-Raphaël]

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[Bourgain]

L^2 blowups **parabolically** concentrate some mass.

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- Extensions in [Merle-Vega], [Carles-Keraani], [Bégout-Vargas].

TYPICAL BLOWUPS LEAVE AN L^2 STAIN AT TIME T^*

[Merle-Raphaël]:

$H^1 \cap \{\|Q\|_{L^2} < \|u_0\|_{L^2} < \|Q\|_{L^2} + \alpha^*\} \ni u_0 \mapsto u$ solving $NLS_3^-(\mathbb{R}^2)$ on $[0, T^*)$ (maximal) with $T^* < \infty$.

$\exists \lambda(t), x(t), \theta(t) \in \mathbb{R}^+, \mathbb{R}^2, \mathbb{R}/(2\pi\mathbb{Z})$ and u^* such that

$$u(t) - \lambda(t)^{-1} Q\left(\frac{x - x(t)}{\lambda(t)}\right) e^{i\theta(t)} \rightarrow u^*$$

strongly in $L^2(\mathbb{R}^2)$. Typically, $u^* \notin H^s \cup L^p$ for $s > 0, p > 2$!

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- **Scattering Below the Ground State Mass**

$$\|u_0\|_{L^2} < \|Q\|_{L^2} \implies ??? \ u_0 \longmapsto u \text{ with } \|u\|_{L_{tx}^4} < \infty.$$

(Also, L^2 solutions of $NLS_3^+(\mathbb{R}^2)$ satisfy^{???} $\|u\|_{L_{tx}^4} < \infty$.)

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■ Minimal Mass Blowup Characterization

$$\|u_0\|_{L^2} = \|Q\|_{L^2}, u_0 \longmapsto u, T^* < \infty \implies ??? \ u = S,$$

modulo a solution symmetry. An intermediate step would extend characterization of the minimal mass blowup solutions in H^s for $s < 1$.

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- Same results for $NLS_{\frac{4}{d}+1}^-(\mathbb{R}^d)$ in H^s , $s > \frac{d+8}{d+10}$. [Visan-Zhang]

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- Spacetime norm divergence rate

$$\|u\|_{L^4_{tx}([0,t]\times\mathbb{R}^2)} \gtrsim (T^* - t)^{-\beta}$$

is linked with mass concentration rate

$$\limsup_{t \uparrow T^*} \sup_{\text{cubes } I, \text{side}(I) \leq (T^* - t)^{\frac{1}{2} + \beta}} \int_I |u(t, x)|^2 dx \geq \|u_0\|_{L^2}^{-M}.$$

[Work in progress with Roudenko]

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- $H^1 \cap \{\text{radial}\} \ni u_0 \longmapsto u, T^* < \infty$ then for any $a > 0$

$$\|\nabla u(t)\|_{L_{|x|<a}^2} \uparrow \infty \text{ as } t \uparrow T^*.$$

Thus, radial solutions must explode at the origin.

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$$\|\nabla u(t)\|_{L^2}^2 \lesssim C(a, Mass[u_0], H[u_0]) + C(a, Mass[u_0])\|\nabla u(t)\|_{L^2_{|x|<a}}^3.$$



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Induction on Energy; Interaction Morawetz; Mass Freezing
- Focusing energy critical case?
 - [Kenig-Merle]: $E[u_0] < E[Q]$ and $\|\nabla u_0\|_{L^2} < \|\nabla Q\|_{L^2} \implies$ global-in-time and scatters.

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- **Conjecture:** $NLS_7^+(\mathbb{R}^3)$ is GWP and scatters in $H^{7/6}(\mathbb{R}^3)$.
[See discussion by Bourgain, GAFA Special Volume, 2000]

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- Frequency heuristic: Bounded $H^{1/2}$ blowup inconsistent with mass conservation.

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The proof follows [Merle-Tsutsumi] with the **(CH)** upper bound as a proxy for L^2 conservation. Explicit constant from sharp Gagliardo-Nirenberg estimate. [Delpino-Dolbeaut]

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■ Fix a number K by the condition

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- **Kth Doubling Time after t_0 :**

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$$\begin{aligned}\|P_{|\xi|>R_M(t_0)}u(t_1)\|_{L^2} &\geq \frac{1}{2}\|P_{|\xi|>R_M}(t_0)u(t_0)\|_{L^2} \\ &\gtrsim \mu_0 R_{M-\mu_0}^{-1/2}(t_0).\end{aligned}$$

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- For small γ_0 , we can not park this mass inside the gap $R_M(t_0) < |\xi| < R_M(t_1)$ so we have to put it in the high frequency boondox $|\xi| > R_M(t_1)$.
- Since, at time t_1 , $R_M(t_1) \geq KR_{M-\mu_0}(t_0)$, we conclude

$$\|u(t_1)\|_{H^{1/2}} \geq 3\Lambda,$$

a contradiction.

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- Assuming that $\|u\|_{L^5_{tx}([0, t] \times \mathbb{R}^3)} \lesssim (T^* - t)^{-1/5+}$ and the **Concentration Property** we **contradict (CH)** proving critical norm explosion.

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- Extends to the general mass supercritical case?